# The Gauss curvature flow : Regularity and Asymptotic Behavior 

Kyeongsu Choi

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# ABSTRACT <br> The Gauss curvature flow : Regularity and Asymptotic Behavior Kyeongsu Choi 

This thesis contains the author's results on the evolution of convex hypersurfaces by positive powers of the Gauss curvature

$$
\frac{\partial}{\partial t} F=K^{\alpha} \vec{n} .
$$

We first establish interior estimates for strictly convex solutions by deriving lower bounds for the principal curvatures and upper bounds for the Gauss curvature. We also investigate the optimal regularity of weakly convex translating solutions. The interesting case is when the translator has flat sides. We prove the existence of such translators and show that they are of optimal class $C^{1,1}$. Finally, we classify all closed self-similar solutions of the Gauss curvature flow which is closely related to the asymptotic behavior.

## Table of contents

CHAPTER 1. Introduction ..... 1
CHAPTER 2. Preliminaries ..... 3
Notation ..... 3
Evolution equations ..... 4
Equations of Self-Shrinkers ..... 7
Euler's formula ..... 8
CHAPTER 3. Regularity of strictly convex solutions ..... 11
Lower bounds on the principal curvatures ..... 12
Speed estimates ..... 16
CHAPTER 4. Optimal regularity of translators ..... 22
Optimal $C^{1,1}$ regularity ..... 24
Partial derivative bound ..... 27
Bistance from the tip to flat sides ..... 34
CHAPTER 5. Asymptotic Behavior of closed solutions ..... 39
Pogorelov type computation ..... 40
Strong maximum principle ..... 43
Bibliography ..... 46

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## Chapter 1

## Introduction

We consider a family of complete non-compact strictly convex hypersurfaces $\Sigma_{t}$ embedded in $\mathbb{R}^{n+1}$ which evolve by any positive power of their Gauss curvature $K$.

Given a complete and strict convex hypersurface $\Sigma_{0}$ embedded in $\mathbb{R}^{n+1}$, we let $F_{0}: M^{n} \rightarrow$ $\mathbb{R}^{n+1}$ be an immersion with $F_{0}\left(M^{n}\right)=\Sigma_{0}$. For a given number $\alpha>0$, we say that an oneparameter family of immersions $F: M^{n} \times(0, T) \rightarrow \mathbb{R}^{n+1}$ is a solution of the $\alpha$-Gauss curvature flow, if for each $t \in(0, T), F\left(M^{n}, t\right)=\Sigma_{t}$ is a complete and strictly convex hypersurface embedded in $\mathbb{R}^{n+1}$, and $F(\cdot, t)$ satisfies

$$
\begin{cases}\frac{\partial}{\partial t} F(p, t) & =K^{\alpha}(p, t) \vec{n}(p, t) \\ \lim _{t \rightarrow 0} F(p, t) & =F_{0}(p)\end{cases}
$$

where $K(p, t)$ is the Gauss curvature of $\Sigma_{t}$ at $F(p, t)$ and $\vec{n}(p, t)$ is the interior unit normal vector of $\Sigma_{t}$ at the point $F(p, t)$.

The Gauss curvature flow was first introduced by W. Firey in [18], where he showed that a closed strictly convex and centrally symmetric solution in $\mathbb{R}^{3}$ converges to a round sphere after rescailing the solution. In [25] K. Tso established the existence of closed and strictly convex solutions in $\mathbb{R}^{n+1}$ and showed that it converges to a point. B. Chow [12] extended Tso's result to the flow by positive powers of the Gauss curvature, namely a strictly convex closed solution to the $\alpha$-Gauss curvature flow $\partial_{t} F=K^{\alpha} \vec{n}$. Schnürer [24] showed the existence of strictly convex entire graph solution to the $\alpha$-Gauss curvature flow for $0<\alpha<\frac{1}{n-1}$. The author, jointly with P. Daskalolpoulos, L. Kim, and K. Lee [10], established the all-times existence of non-compact, complete, and strictly convex solution solutions to the $\alpha$-Gauss curvature flow for $\alpha>0$. In addition, J. Urbas [26], [27] showed that given an exponent $\alpha>\frac{1}{2}$ and a strictly convex open bounded domain $\Omega \subset \mathbb{R}^{n+1}$, there exists a complete smooth strictly convex graph translating by the $\alpha$-power of the Gauss curvature $K^{\alpha}$.

Regarding weakly convex solutions, B. Andrews [1] showed the optimal $C^{1,1}$ regularity of the solution to the Gauss curvature flow in the case $n=2$, and extended to the result to $\alpha \in\left[\frac{1}{2}, 1\right]$ and $n=2$ with X. Chen [3]. In [16] the $C^{1, \beta}$ regularity of weakly convex viscosity solutions to the $\alpha$-Gauss curvature flow was established. The author, jointly with P. Daskalolpoulos and K. Lee [11], established the optimal $C^{1,1}$ regularity to the translating solutions to the Gauss curvature flow in the case $n=2$. They also showed that the translator defined on a square has flat sides so that $C^{1,1}$ is actually the optimal regularity.

In particular, if the initial closed convex hypersurface $\Sigma_{0}$ has a flat side, then the solution $\Sigma_{t}$ to the $\alpha$-Gauss curvature flow with $\alpha>\frac{1}{n}$ preserves the flat side for some finite time. See [14], [15], [23]. If $\alpha \leqslant \frac{1}{n}$, B. Andrews [2] showed that any weakly convex closed hypersurface becomes strictly convex immediately by the $\alpha$-Gauss curvature flow.

The convergence of the flow to a closed self-similar solution has been widely studied in $[\mathbf{1}, \mathbf{2}, \mathbf{3}, \mathbf{4}, \mathbf{1 2}, \mathbf{1 3}, \mathbf{1 9}, \mathbf{2 0}, 21,22]$ for $\alpha>\frac{1}{n+2}$. In the case $\alpha=\frac{1}{n+2}$, E. Calabi [7] that showed that closed self-similar solutions are ellipsoids. In the case $\alpha=\frac{1}{n}$, B. Chow [12] established the convergence to the round sphere, namely the sphere is the unique closed self-similar solution. In the two dimension case $n=2$, B. Andrews showed the convergence to sphere for $\alpha=1$ in [1], and for $\alpha \in\left(\frac{1}{2}, 1\right)$ with X. Chen in [3]. P. Guan and L. Ni [19] obtained the convergence of the Gauss curvature flow to a self-similar solution and in [4] they extended the same result to $\alpha \geqslant \frac{1}{n+2}$ jointly with B. Andrews. So, it remained to classify the closed self-similar solutions $\Sigma_{t}=(T-t)^{\frac{1}{1+n a}} \Sigma$. Namely, it is enough to classify strictly convex closed hypersurfaces $\Sigma$ satisfying

$$
K^{\alpha}(p)=-\langle F(p), \vec{n}(p)\rangle .
$$

The author and P. Daskalopoulus [9] showed the uniqueness of the closed self-similar solution to the $\alpha$-Gauss curvature flow for $\alpha \in\left[\frac{1}{n}, 1+\frac{1}{n}\right.$ ), and jointly with S . Brendle, they extended the result to $\alpha>\frac{1}{n+2}$. Moreover, in the case $\alpha=\frac{1}{n+2}$, they gave an alternating proof of the classification result of Calabi [7].

## Chapter 2

## Preliminaries

## 1 Notation

We summarize the following notation which will be frequently used in what follows.
(i) We recall that $g_{i j}=\left\langle F_{i}, F_{j}\right\rangle$, where $F_{i}:=\nabla_{i} F$. Also, we denote as usual by $g^{i j}$ the inverse matrix of $g_{i j}$ and $F^{i}=g^{i j} F_{j}$.
(ii) Assume $\Sigma_{t}$ is a strictly convex graph solution of $\sqrt{\left.*^{\alpha}\right\}}$ in $\mathbb{R}^{n+1}$. Then, we let $\bar{u}(\cdot, t)$ : $M^{n} \rightarrow \mathbb{R}$ denote the height function $\bar{u}(p, t)=\left\langle F(p, t), \vec{e}_{n+1}\right\rangle$.
(iii) Given constants $M$ and $\beta \geqslant 0$, we define the cut-off function $\psi_{\beta}$ by

$$
\psi_{\beta}(p, t)=(M-\beta t-\bar{u}(p, t))_{+}=\max \{M-\beta t-\bar{u}(p, t), 0\} .
$$

In particular, we denote $\psi_{0}:=(M-\bar{u}(p, t))_{+}$by $\psi$ for convenience.
Also, given a constant $R>0$ and a point $Y \in \mathbb{R}^{n+1}$, we define the cut-off function $\bar{\psi}$ by

$$
\bar{\psi}(p, t)=\left(R^{2}-|F(p, t)-Y|^{2}\right)_{+}=\max \left\{R^{2}-|F(p, t)-Y|^{2}, 0\right\}
$$

(iv) Given a ball $B_{R}^{n+1}(Y) \subset \mathbb{R}^{n+1}$ and a complete hypersurface $\Sigma \subset \mathbb{R}^{n+1}$, we say that a compact hypersurface $\Sigma^{c}$ with boundary $\partial \Sigma^{c}$ is cut off from $\Sigma$ by $B_{R}^{n+1}(Y)$, if $\Sigma^{c} \subset \Sigma$ and $\partial \Sigma^{c} \subset \partial B_{R}^{n+1}(Y)$ hold. Moreover, Given an immersion $F: M^{n} \times[0, T) \rightarrow \mathbb{R}^{n+1}$ defining $F\left(M^{n}, t\right)=\Sigma_{t}$ and a ball $B_{R}^{n+1}(Y)$ cutting $\Sigma_{t}$, we define a cut-off function $\eta: M^{n} \times[0, T) \rightarrow \mathbb{R}^{n+1}$ by

$$
\eta(p, t)=\left(|F(p, t)-Y|^{2}-R^{2}\right)_{+} .
$$

(v) For a strictly convex smooth hypersurface $\Sigma_{t}$, we denote by $b^{i j}$ the the inverse matrix $\left(h^{-1}\right)^{i j}$ of its second fundamental form $h_{i j}$, namely $b^{i j} h_{j k}=\delta_{k}^{i}$. Notice that the eigenvalues of $b^{i j}$ on an orthonormal frame are the principal radii of curvature.


Figure 1. Cutting ball
(vi) We denote by $\mathcal{L}$ the linearized operator

$$
\mathcal{L}=\alpha K^{\alpha} b^{i j} \nabla_{i} \nabla_{j}
$$

Furthermore, $\langle,\rangle_{\mathcal{L}}$ denotes the associated inner product $\langle\nabla f, \nabla g\rangle_{\mathcal{L}}=\alpha K^{\alpha} b^{i j} \nabla_{i} f \nabla_{j} g$, where $f, g$ are differentiable functions on $M^{n}$, and $\|\cdot\|_{\mathcal{L}}$ denotes the $\mathcal{L}$-norm given by the inner product $\langle,\rangle_{\mathcal{L}}$
(vii) We denote by $v=\left\langle\vec{n}, \vec{e}_{n+1}\right\rangle^{-1}$ the gradient function.
(viii) $H$ and $\lambda_{\min }$ denote the mean curvature and the smallest principal curvature, respectively.
(ix) We will use in the sequel the functions $Z: M^{n} \rightarrow \mathbb{R}$ and $W: M^{n} \rightarrow \mathbb{R}$ defined by

$$
Z(p)=\left(K^{\alpha} b^{i j} g_{i j}-\frac{n \alpha-1}{2 \alpha}|F|^{2}\right)(p), \quad W(p)=\left(K^{\alpha} \lambda_{\min }^{-1}-\frac{n \alpha-1}{2 n \alpha}|F|^{2}\right)(p)
$$

## 2 Evolution equations

In this section we will derive some basic equations of the $\alpha$-Gauss curvature flow.
Proposition 2.1. Assume that $\Sigma_{t}$ is a complete strictly convex graph solution of $\left(*^{\alpha}\right)$. Then, the following hold
(2.1) $\partial_{t} g_{i j}=-2 K^{\alpha} h_{i j}$
(2.2) $\partial_{t} g^{i j}=2 K^{\alpha} h^{i j}$
(2.3) $\partial_{t} \vec{n}=-\nabla K^{\alpha}:=-\left(\nabla_{j} K^{\alpha}\right) F^{j}$
(2.4) $\partial_{t} \psi_{\beta}=\mathcal{L} \psi+(n \alpha-1) v^{-1} K^{\alpha}-\beta$
(2.5) $\partial_{t} \bar{\psi} \leqslant \mathcal{L} \bar{\psi}+2(n \alpha+1)\left(\lambda_{\min }^{-1}+R\right) K^{\alpha}$
(2.6) $\partial_{t} h_{i j}=\mathcal{L} h_{i j}+\alpha K^{\alpha}\left(\alpha b^{k l} b^{m n}-b^{k m} b^{n l}\right) \nabla_{i} h_{m n} \nabla_{j} h_{k l}+\alpha K^{\alpha} H h_{i j}-(1+n \alpha) K^{\alpha} h_{i k} h_{j}^{k}$
(2.7) $\partial_{t} K^{\alpha}=\mathcal{L} K^{\alpha}+\alpha K^{2 \alpha} H$
(2.8) $\quad \partial_{t} b^{p q}=\mathcal{L} b^{p q}-\alpha K^{\alpha} b^{i p} b^{j q}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right) \nabla_{i} h_{k l} \nabla_{j} h_{m n}-\alpha K^{\alpha} H b^{p q}+(1+n \alpha) K^{\alpha} g^{p q}$
(2.9) $\partial_{t} v=\mathcal{L} v-2 v^{-1}\|\nabla v\|_{\mathcal{L}}^{2}-\alpha K^{\alpha} H v$

Proof of 2.1. Observe $\partial_{t} g_{i j}=\left\langle\partial_{t} \nabla_{i} F, \nabla_{j} F\right\rangle+\left\langle\nabla_{i} F, \partial_{t} \nabla_{j} F\right\rangle=\left\langle\nabla_{i} \partial_{t} F, \nabla_{j} F\right\rangle+\left\langle\nabla_{i} F, \nabla_{j} \partial_{t} F\right\rangle$. Hence, $\left\langle\nabla_{i} F, \partial_{t} F\right\rangle=0$ gives $\partial_{t} g_{i j}=-2\left\langle\partial_{t} F, \nabla_{i} \nabla_{j} F\right\rangle=-2\left\langle K^{\alpha} \vec{n}, h_{i j} \vec{n}\right\rangle=-2 K^{\alpha} h_{i j}$.

Proof of 2.2. From $g^{i j} g_{j k}=\delta_{k}^{i}$, we can derive $\partial_{t} g^{i j}=-g^{i k} g^{j l} \partial_{t} g_{k l}=2 K^{\alpha} g^{i k} g^{j l} h_{k l}=2 K^{\alpha} h^{i j}$.
Proof of $(2.3) .|\vec{n}|^{2}=1$ implies $\left\langle\partial_{t} \vec{n}, \vec{n}\right\rangle=0$. Also, $\left\langle\vec{n}, \nabla_{i} F\right\rangle=0$ leads to

$$
\left\langle\partial_{t} \vec{n}, \nabla_{i} F\right\rangle=-\left\langle\vec{n}, \partial_{t} \nabla_{i} F\right\rangle=-\left\langle\vec{n}, \nabla_{i} \partial_{t} F\right\rangle=-\left\langle\vec{n}, \nabla_{i}\left(K^{\alpha} \vec{n}\right)\right\rangle=-\nabla_{i} K^{\alpha}
$$

from which (2.3) readily follows.
Proof of 2.4. The definition of the linearized operator $\mathcal{L}=: \alpha K^{\alpha} b^{i j} \nabla_{i} \nabla_{j}$ gives

$$
\begin{equation*}
\mathcal{L} F:=\alpha K^{\alpha} b^{i j} \nabla_{i} \nabla_{j} F=\alpha K^{\alpha} b^{i j} h_{i j} \vec{n}=n \alpha K^{\alpha} \vec{n}=(n \alpha) \partial_{t} F \tag{2.10}
\end{equation*}
$$

which yields that $\mathcal{L} \bar{u}=(n \alpha) \partial_{t} \bar{u}$. Therefore,

$$
\partial_{t} \psi_{\beta}=-\beta-\partial_{t} \bar{u}=-\mathcal{L} \bar{u}+(n \alpha-1) \partial_{t} \bar{u}-\beta
$$

holds on the support of $\psi_{\beta}:=(M-\beta t-\bar{u}(p, t))_{+}$. Substituting for $\partial_{t} \bar{u}=\left\langle\partial_{t} F, \vec{e}_{n+1}\right\rangle=$ $\left\langle K^{\alpha} \vec{n}, \vec{e}_{n+1}\right\rangle=K^{\alpha} v^{-1}$, where $v:=\left\langle\vec{n}, \vec{e}_{n+1}\right\rangle^{-1}$, yields (2.4).

Proof of (2.5). By (2.10), on the support of $\bar{\psi}=:\left(R^{2}-|F-Y|^{2}\right)_{+}$we have

$$
\begin{aligned}
\partial_{t} \bar{\psi} & =-2\left\langle\left(\mathcal{L} F-(n \alpha) \partial_{t} F\right)+\partial_{t} F, F-Y\right\rangle=-2\langle\mathcal{L} F, F-Y\rangle+2(n \alpha-1)\left\langle K^{\alpha} \vec{n}, F-Y\right\rangle \\
& \leqslant \mathcal{L} \bar{\psi}+2\|\nabla F\|_{\mathcal{L}}^{2}+2 K^{\alpha}|n \alpha-1||F-Y| \leqslant \mathcal{L} \bar{\psi}+2 \alpha K^{\alpha} b^{i j} g_{i j}+2(n \alpha+1) R K^{\alpha} \\
& \leqslant \mathcal{L} \bar{\psi}+2 n \alpha \lambda_{\min }^{-1} K^{\alpha}+2(n \alpha+1) R K^{\alpha} \leqslant \mathcal{L} \bar{\psi}+2(n \alpha+1)\left(\lambda_{\min }^{-1}+R\right) K^{\alpha} .
\end{aligned}
$$

Proof of (2.6). We have

$$
\begin{aligned}
\partial_{t} h_{i j} & =\partial_{t}\left\langle\nabla_{i} \nabla_{j} F, \vec{n}\right\rangle=\left\langle\nabla_{i} \nabla_{j} \partial_{t} F, \vec{n}\right\rangle+\left\langle\nabla_{i} \nabla_{j} F, \partial_{t} \vec{n}\right\rangle=\left\langle\nabla_{i} \nabla_{j}\left(K^{\alpha} \vec{n}\right), \vec{n}\right\rangle+\left\langle h_{i j} \vec{n}, \partial_{t} \vec{n}\right\rangle \\
& =\nabla_{i} \nabla_{j} K^{\alpha}+\left\langle\nabla_{j} K^{\alpha} \nabla_{i} \vec{n}, \vec{n}\right\rangle+\left\langle\nabla_{i} K^{\alpha} \nabla_{j} \vec{n}, \vec{n}\right\rangle+K^{\alpha}\left\langle\nabla_{i} \nabla_{j} \vec{n}, \vec{n}\right\rangle+0 .
\end{aligned}
$$

By $\left\langle\nabla_{i} \vec{n}, \vec{n}\right\rangle=0$ and $\left\langle\nabla_{i} \nabla_{j} \vec{n}, \vec{n}\right\rangle=-\left\langle\nabla_{j} \vec{n}, \nabla_{i} \vec{n}\right\rangle=-h_{i m} h_{j}^{m}$, we deduce

$$
\begin{equation*}
\partial_{t} h_{i j}=\nabla_{i} \nabla_{j} K^{\alpha}-K^{\alpha} h_{i m} h_{j}^{m} . \tag{2.11}
\end{equation*}
$$

Applying $K=\operatorname{det}\left(h_{i j} g^{j k}\right)=\operatorname{det}\left(h_{i j}\right) \operatorname{det}\left(g^{k l}\right), \nabla \log \operatorname{det}\left(h_{i j}\right)=b^{k l} \nabla h_{k l}$ and $\partial b^{k l} / \partial h_{m n}=-b^{k m} b^{l n}$ on the first term on the right hand side of (2.11), yields

$$
\begin{aligned}
\partial_{t} h_{i j} & =\nabla_{i} \nabla_{j} K^{\alpha}-K^{\alpha} h_{i m} h_{j}^{m}=\nabla_{i}\left(\alpha K^{\alpha} b^{k l} \nabla_{j} h_{k l}\right)-K^{\alpha} h_{i m} h_{j}^{m} \\
& =\alpha^{2} K^{\alpha} b^{k l} b^{m n} \nabla_{i} h_{m n} \nabla_{j} h_{k l}+\alpha K^{\alpha} \frac{\partial b^{k l}}{\partial h_{m n}} \nabla_{i} h_{m n} \nabla_{j} h_{k l}+\alpha K^{\alpha} b^{k l} \nabla_{i} \nabla_{j} h_{k l}-K^{\alpha} h_{i m} h_{j}^{m} \\
& =\alpha K^{\alpha} b^{k l} \nabla_{i} \nabla_{j} h_{k l}+\alpha K^{\alpha}\left(\alpha b^{k l} b^{m n}-b^{k m} b^{l n}\right) \nabla_{i} h_{m n} \nabla_{j} h_{k l}-K^{\alpha} h_{i m} h_{j}^{m} .
\end{aligned}
$$

On the other hand,

$$
\begin{aligned}
\alpha K^{\alpha} b^{k l} \nabla_{i} \nabla_{j} h_{k l} & =\alpha K^{\alpha} b^{k l} \nabla_{i} \nabla_{k} h_{j l}=\alpha K^{\alpha} b^{k l} \nabla_{k} \nabla_{i} h_{j l}+\alpha K^{\alpha} b^{k l} R_{i k j m} h_{l}^{m}+\alpha K^{\alpha} b^{k l} R_{i k l m} h_{j}^{m} \\
& =\alpha K^{\alpha} b^{k l} \nabla_{k} \nabla_{l} h_{i j}+\alpha K^{\alpha}\left(h_{i j} h_{k m}-h_{i m} h_{k j}\right) h_{l}^{m} b^{l k}+\alpha K^{\alpha}\left(h_{i l} h_{k m} b^{k l}-h_{i m} h_{k l} b^{k l}\right) h_{j}^{m} \\
& =\mathcal{L} h_{i j}+\alpha K^{\alpha}\left(h_{i j} h_{k m}-h_{i m} h_{k j}\right) g^{m k}+\alpha K^{\alpha}\left(h_{i m}-n h_{i m}\right) h_{j}^{m} \\
& =\mathcal{L} h_{i j}+\alpha K^{\alpha} H h_{i j}-n \alpha K^{\alpha} h_{i m} h_{j}^{m}
\end{aligned}
$$

and (2.6) easily follows.

Proof of (2.7). From (2.11) we have

$$
\mathcal{L} K^{\alpha}=\alpha K^{\alpha} b^{i j}\left(\partial_{t} h_{i j}+K^{\alpha} h_{i m} h_{j}^{m}\right)=\alpha K^{\alpha} b^{i j} \partial_{t} h_{i j}+\alpha K^{2 \alpha} H .
$$

Also, from $K=\operatorname{det}\left(h_{i j}\right) \operatorname{det}\left(g^{k l}\right)$, we derive that

$$
\begin{aligned}
\partial_{t} K^{\alpha} & =\alpha K^{\alpha} \partial_{t}\left(\log \left(\operatorname{det} h_{i j}\right)+\log \left(\operatorname{det} g^{k l}\right)\right)=\alpha K^{\alpha} b^{i j} \partial_{t} h_{i j}+\alpha K^{\alpha} g_{k l} \partial_{t} g^{k l} \\
& =\mathcal{L} K^{\alpha}-\alpha K^{2 \alpha} H+\alpha K^{\alpha} g_{k l} \partial_{t} g^{k l} .
\end{aligned}
$$

Aplying (2.2) on the last term yields (2.7).

Proof of 2.8 . The identity $b^{i j} h_{j k}=\delta_{k}^{i}$ leads to

$$
\begin{align*}
\partial_{t} b^{p q} & =-b^{i p} b^{j q} \partial_{t} h_{i j}  \tag{2.12}\\
\nabla_{m} b^{p q} & =-b^{i p} b^{j q} \nabla_{m} h_{i j} \tag{2.13}
\end{align*}
$$

Therefore,

$$
\begin{aligned}
\nabla_{n} \nabla_{m} b^{p q} & =-b^{j q} \nabla_{n} b^{i p} \nabla_{m} h_{i j}-b^{i p} \nabla_{n} b^{j q} \nabla_{m} h_{i j}-b^{i p} b^{j q} \nabla_{n} \nabla_{m} h_{i j} \\
& =2 b^{j q} b^{i k} b^{p l} \nabla_{n} h_{k l} \nabla_{m} h_{i j}-b^{i p} b^{j q} \nabla_{n} \nabla_{m} h_{i j} .
\end{aligned}
$$

Hence, $\mathcal{L} b^{p q}:=\alpha K^{\alpha} b^{n m} \nabla_{n} \nabla_{m} b^{p q}$ satisfies

$$
\mathcal{L} b^{p q}=2 \alpha K^{\alpha} b^{n m} b^{j q} b^{i k} b^{p l} \nabla_{n} h_{k l} \nabla_{m} h_{i j}-b^{i p} b^{j q} \mathcal{L} h_{i j}=2 \alpha K^{\alpha} b^{l p} b^{j q} b^{k i} b^{n m} \nabla_{l} h_{k n} \nabla_{j} h_{i m}-b^{i p} b^{j q} \mathcal{L} h_{i j}
$$

Combing the last identity with (2.6) and (2.12) yields

$$
\begin{aligned}
\partial_{t} b^{p q} & =-b^{i p} b^{j q}\left(\mathcal{L} h_{i j}+\alpha K^{\alpha}\left(\alpha b^{k l} b^{m n}-b^{k m} b^{n l}\right) \nabla_{i} h_{k l} \nabla_{j} h_{m n}+\alpha K^{\alpha} H h_{i j}-(1+n \alpha) K^{\alpha} h_{i k} h_{j}^{k}\right) \\
& =\mathcal{L} b^{p q}-\alpha K^{\alpha} b^{i p} b^{j q}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right) \nabla_{i} h_{k l} \nabla_{j} h_{m n}-\alpha K^{\alpha} H b^{p q}+(1+n \alpha) K^{\alpha} g^{p q}
\end{aligned}
$$

which gives (2.8).

Proof of 2.9. By 2.3) we have $\partial_{t} v=-v^{2}\left\langle\vec{e}_{n+1}, \partial_{t} \vec{n}\right\rangle=v^{2}\left\langle\vec{e}_{n+1}, \nabla K^{\alpha}\right\rangle$. Furthermore,

$$
\begin{aligned}
\mathcal{L} v & =-\alpha K^{\alpha} b^{i j} \nabla_{i}\left(v^{2}\left\langle\vec{e}_{n+1}, \nabla_{j} \vec{n}\right\rangle\right)=-2 \alpha K^{\alpha} b^{i j} v \nabla_{i} v\left\langle\vec{e}_{n+1}, \nabla_{j} \vec{n}\right\rangle+\alpha K^{\alpha} b^{i j} v^{2}\left\langle\vec{e}_{n+1}, \nabla_{i}\left(h_{j k} F^{k}\right)\right\rangle \\
& =2 v^{-1}\|\nabla v\|_{\mathcal{L}}^{2}+\alpha K^{\alpha} b^{i j} h_{j k} h_{i}^{k} v^{2}\left\langle\vec{e}_{n+1}, \vec{n}\right\rangle+v^{2}\left\langle\vec{e}_{n+1}, \alpha K^{\alpha} b^{i j}\left(\nabla_{i} h_{j k}\right) F^{k}\right\rangle \\
& =2 v^{-1}\|\nabla v\|_{\mathcal{L}}^{2}+\alpha K^{\alpha} H v+v^{2}\left\langle\vec{e}_{n+1}, \nabla K^{\alpha}\right\rangle
\end{aligned}
$$

Combining the above yields (2.9).

## 3 Equations of Self-Shrinkers

In this section we will derive basic equations of the closed self-similar solutions $\Sigma$ to the $\alpha$ Gauss curvature flow. Since $(T-t)^{\frac{1}{n \alpha+1}} \Sigma$ is a solution to the $\alpha$-Gauss curvature flow, an immersion $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ of $F\left(M^{n}\right)=\Sigma$ satisfies $\left(*^{\alpha}\right)$

Proposition 2.2. Given a strictly convex smooth solution $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ to $\left(*^{\alpha}\right)$, the following hold

$$
\begin{align*}
\nabla_{i} b^{j k}= & -b^{j l} b^{k m} \nabla_{i} h_{l m},  \tag{2.14}\\
\mathcal{L}|F|^{2}= & 2 \alpha K^{\alpha} b^{i j}\left(g_{i j}-h_{i j} K^{\alpha}\right)=2 \alpha K^{\alpha} b^{i j} g_{i j}-2 n \alpha K^{2 \alpha},  \tag{2.15}\\
\nabla_{i} K^{\alpha}= & h_{i j}\left\langle F, F^{j}\right\rangle,  \tag{2.16}\\
\mathcal{L} K^{\alpha}= & \left\langle F, \nabla K^{\alpha}\right\rangle+n \alpha K^{\alpha}-\alpha K^{2 \alpha} H,  \tag{2.17}\\
\mathcal{L} b^{p q}= & K^{-\alpha} b^{p r} b^{q s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}+\alpha K^{\alpha} b^{p r} b^{q s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m}  \tag{2.18}\\
& +\left\langle F, \nabla b^{p q}\right\rangle-b^{p q}-(n \alpha-1) g^{p q} K^{\alpha}+\alpha K^{\alpha} H b^{p q} .
\end{align*}
$$

Proof. From $\nabla_{i}\left(b^{j k} h_{k l}\right)=\nabla_{i} \delta_{l}^{j}=0$, we can derive $h_{k l} \nabla_{i} b^{j k}=-b^{j k} \nabla_{i} h_{k l}$. Hence, we have (2.14) by

$$
\nabla_{i} b^{j m}=b^{l m} h_{k l} \nabla_{i} b^{j k}=-b^{l m} b^{j k} \nabla_{i} h_{k l} .
$$

Also, by definition $\mathcal{L}:=\alpha K^{\alpha} b^{i j} \nabla_{i} \nabla_{j}$ we have

$$
\mathcal{L}|F|^{2}=2 \alpha K^{\alpha} b^{i j}\left\langle F_{i}, F_{j}\right\rangle+2 \alpha K^{\alpha} b^{i j}\left\langle F, \nabla_{i} \nabla_{j} F \vec{n}\right\rangle=2 \alpha K^{\alpha} b^{i j} g_{i j}+2 \alpha K^{\alpha} b^{i j}\left\langle F, h_{i j} \vec{n}\right\rangle .
$$

Thus, the given equation $\left(\sqrt{\left.* *^{\alpha}\right)}\right.$ implies (5.5).
Equation (2.16) can be simply obtained by differentiating (***

$$
\nabla_{i} K^{\alpha}=h_{i k}\left\langle F, F^{k}\right\rangle .
$$

Differentiating the equation above again we obtain

$$
\nabla_{i} \nabla_{j} K^{\alpha}=\nabla_{i} h_{j k}\left\langle F, F^{k}\right\rangle+h_{i j}+h_{i k} h_{j}^{k}\langle F, \vec{n}\rangle=\left\langle F, \nabla h_{i j}\right\rangle+h_{i j}-h_{i k} h_{j}^{k} K^{\alpha} .
$$

On the other hand, (2.14) and direct differentiation yield
$\nabla_{i} \nabla_{j} K^{\alpha}=\nabla_{i}\left(\alpha K^{\alpha} b^{p q} \nabla_{j} h_{p q}\right)=\alpha K^{\alpha} b^{p q} \nabla_{i} \nabla_{j} h_{p q}+\alpha^{2} K^{\alpha} b^{r s} b^{p q} \nabla_{i} h_{r s} \nabla_{j} h_{p q}-\alpha K^{\alpha} b^{p r} b^{q s} \nabla_{i} h_{r s} \nabla_{j} h_{p q}$.
Observing

$$
\begin{aligned}
\nabla_{i} \nabla_{j} h_{p q}=\nabla_{i} \nabla_{p} h_{j q} & =\nabla_{p} \nabla_{i} h_{j q}+R_{i p j m} h_{q}^{m}+R_{i p q m} h_{j}^{m} \\
& =\nabla_{p} \nabla_{q} h_{i j}+\left(h_{i j} h_{p m}-h_{i m} h_{j p}\right) h_{q}^{m}+\left(h_{i q} h_{p m}-h_{i m} h_{p q}\right) h_{j}^{m}
\end{aligned}
$$

we obtain
$\alpha K^{\alpha} b^{p q} \nabla_{i} \nabla_{j} h_{p q}=\alpha K^{\alpha} b^{p q} \nabla_{p} \nabla_{q} h_{i j}+\alpha K^{\alpha} H h_{i j}-n \alpha K^{\alpha} h_{i m} h_{j}^{m}=\mathcal{L} h_{i j}+\alpha K^{\alpha} H h_{i j}-n \alpha K^{\alpha} h_{i m} h_{j}^{m}$.
Combining the equations above yields

$$
\begin{align*}
\mathcal{L} h_{i j}= & -\alpha^{2} K^{\alpha} b^{r s} b^{p q} \nabla_{i} h_{r s} \nabla_{j} h_{p q}+\alpha K^{\alpha} b^{p r} b^{q s} \nabla_{i} h_{r s} \nabla_{j} h_{p q}  \tag{2.6}\\
& +\left\langle F, \nabla h_{i j}\right\rangle+h_{i j}+(n \alpha-1) h_{i k} h_{j}^{k} K^{\alpha}-\alpha K^{\alpha} H h_{i j} .
\end{align*}
$$

We now observe

$$
\begin{aligned}
\mathcal{L} K^{\alpha} & =\alpha K^{\alpha} b^{i j} \nabla_{i}\left(\alpha K^{\alpha} b^{p q} \nabla_{j} h_{p q}\right) \\
& =\alpha^{3} K^{2 \alpha} b^{i j} b^{p q} b^{r s} \nabla_{i} h_{r s} \nabla_{j} h_{p q}-\alpha^{2} K^{2 \alpha} b^{i j} b^{p r} b^{q s} \nabla h_{r s} \nabla_{j} h_{p q}+\alpha K^{\alpha} b^{p q} \mathcal{L} h_{p q}
\end{aligned}
$$

which gives (2.17), since

$$
\mathcal{L} K^{\alpha}=\alpha K^{\alpha} b^{i j}\left(\left\langle F, \nabla h_{i j}\right\rangle+h_{i j}+(n \alpha-1) h_{i k} h_{j}^{k} K^{\alpha}-\alpha K^{\alpha} H h_{i j}\right)=\left\langle F, \nabla K^{\alpha}\right\rangle+n \alpha K^{\alpha}-\alpha K^{2 \alpha} H .
$$

Finally, by using (2.14), we can derive

$$
\mathcal{L} b^{p q}=\alpha K^{\alpha} b^{i j} \nabla_{i}\left(-b^{p r} b^{q s} \nabla_{j} h_{r s}\right)=2 \alpha K^{\alpha} b^{i j} b^{p k} b^{r m} b^{q s} \nabla_{i} h_{k m} \nabla_{j} h_{r s}-b^{p r} b^{q s} \mathcal{L} h_{r s}
$$

Applying (2.6) yields

$$
\begin{aligned}
\mathcal{L} b^{p q}= & \alpha^{2} K^{\alpha} b^{p r} b^{q s} b^{i j} b^{k m} \nabla_{r} h_{i j} \nabla_{s} h_{k m}+\alpha K^{\alpha} b^{p r} b^{q s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m} \\
& +\left\langle F, \nabla b^{p q}\right\rangle-b^{p q}-(n \alpha-1) g^{p q} K^{\alpha}+\alpha K^{\alpha} H b^{p q} .
\end{aligned}
$$

Thus, $\nabla K^{\alpha}=\alpha K^{\alpha} b^{i j} \nabla h_{i j}$ gives the desired result.

## 4 Euler's formula

Finally, we recall Euler's formula to use Pogorelov computations.
Proposition 2.3. Let $\Sigma$ be a smooth hypersurface and $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ be a smooth immersion with $F\left(M^{n}\right)=\Sigma$. Then, for all $p \in M^{n}$ and $i \in\{1, \cdots, n\}$, the following holds

$$
\frac{h_{i i}(p)}{g_{i i}(p)} \leqslant \lambda_{\max }(p)
$$

We can simply modify Euler's formula to $b^{i j}$ as follows:
Proposition 2.4. Let $\Sigma$ be a smooth strictly convex hypersurface and $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ be a smooth immersion with $F\left(M^{n}\right)=\Sigma$. Then, for all $p \in M^{n}$ and $i \in\{1, \cdots, n\}$, the following holds

$$
\frac{b^{i i}(p)}{g^{i i}(p)} \leqslant \frac{1}{\lambda_{\min }(p)}
$$

Proof. Let $\left\{E_{1}, \cdots, E_{n}\right\}$ be an orthonormal basis of $T \Sigma_{F(p)}$ satisfying $L\left(E_{j}\right)=\lambda_{j} E_{j}$, where $L$ is the Weingarten map and $\lambda_{1}, \cdots, \lambda_{n}$ are the principal curvatures of $\Sigma$ at $p$. Denote by $\left\{a_{i j}\right\}$ the matrix satisfying $F_{i}(p):=\nabla_{i} F(p)=a_{i j} E_{j}$ and by $\left\{c_{i j}\right\}$ the diagonal matrix $\operatorname{diag}\left(\lambda_{1}, \cdots, \lambda_{n}\right)$. Then, $L F_{i}(p)=h_{i j}(p) F^{j}(p)$ implies

$$
b^{i j}(p) L F_{j}(p)=b^{i j}(p) h_{j k}(p) F^{k}(p)=F^{i}(p)=g^{i j}(p) a_{j m} E_{m}
$$

Observing that for the sum $a_{j k} E_{k}=\sum_{k=1}^{n} a_{j k} E_{k}$ we have $L\left(a_{j k} E_{k}\right)=a_{j k} L\left(E_{k}\right)=a_{j k} \lambda_{k} E_{k}=$ $a_{j k} c_{k m} E_{m}$, it follows that

$$
g^{i j}(p) a_{j m} E_{m}=b^{i j}(p) L F_{j}(p)=b^{i j}(p) L\left(a_{j k} E_{k}\right)=b^{i j}(p) a_{j k} c_{k m} E_{m}
$$

Hence,

$$
g^{i j}(p) a_{j m}=b^{i j}(p) a_{j k} c_{k m}
$$

Denoting by $a^{i j}$ and $c^{i j}$ the inverse matrices of $a_{i j}$ and $c_{i j}$ respectively, we have

$$
g^{i j}(p) a_{j m} c^{m l} a^{l n}=b^{i j}(p) a_{j k} c_{k m} c^{m l} a^{l n}=b^{i n}(p)
$$

Thus, for each $i \in\{1, \cdots, n\}$, the following holds

$$
b^{i i}(p)=\sum_{j, l, m} g^{i j}(p) a_{j m} c^{m l} a^{l i}=\sum_{j, l, m}\left\langle F^{i}(p), F^{j}(p)\right\rangle a_{j m} c^{m l} a^{l i}
$$

By setting $F^{i}(p)=\bar{a}_{i j} E_{j}$, we observe $a^{j i}=a^{j k}\left\langle F_{k}(p), F^{i}(p)\right\rangle=a^{j k}\left\langle a_{k l} E_{l}, \bar{a}_{i m} E_{m}\right\rangle=a^{j k} a_{k l} \bar{a}_{i m} \delta_{m l}=$ $\bar{a}_{i j}$. Thus, we have $\left\langle F^{i}(p), F^{j}(p)\right\rangle=\left\langle a^{k i} E_{k}, a^{r j} E_{r}\right\rangle=a^{k i} a^{k j}$, which yields

$$
\begin{aligned}
b^{i i}(p) & =\sum_{j, k, l, m} a^{k i} a^{k j} a_{j m} c^{m l} a^{l i}=\sum_{k, l, m} a^{k i} \delta_{m}^{k} c^{m l} a^{l i}=\sum_{k, l} a^{k i} c^{k l} a^{l i}=\sum_{k} a^{k i} a^{k i} \lambda_{k}^{-1}=\sum_{k}\left(a^{k i}\right)^{2} \lambda_{k}^{-1} \\
& \leqslant \sum_{k}\left(a^{k i}\right)^{2} \lambda_{\min }^{-1}=\lambda_{\min }^{-1} \sum_{k, j}\left\langle a^{k i} E_{k}, a^{j i} E_{j}\right\rangle=\lambda_{\min }^{-1}\left\langle F^{i}(p), F^{i}(p)\right\rangle=\lambda_{\min }^{-1} g^{i i}(p)
\end{aligned}
$$

Proposition 2.5. Let $\Sigma$ be a strictly convex smooth hypersurface and $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ be a smooth immersion with $F\left(M^{n}\right)=\Sigma$. Then, for all $p \in M^{n}$ and $i \in\{1, \cdots, n\}$, the following holds

$$
\frac{b^{i 1} g^{i j} b^{j 1}(p)}{g_{11}(p)} \leqslant \frac{1}{\lambda_{\min }^{2}(p)}
$$

Proof. For a fixed point $p \in M^{n}$, we choose an orthonormal basis $\left\{E_{1}, \cdots, E_{n}\right\}$ of $T \Sigma_{F(p)}$ such that $L\left(E_{j}\right)=\lambda_{j} E_{j}$, where $L$ is the Weingarten map and $\lambda_{1}, \cdots, \lambda_{n}$ are the principal curvatures of $\Sigma$ at $p$. Given a chart $(\varphi, U)$ of $p \in \varphi(U) \subset M^{n}$, we denote by $\left\{a_{i j}\right\}$ the matrix satisfying $F_{i}(p):=\nabla_{i} F(p)=a_{i j} E_{j}$ and by $\left\{c_{i j}\right\}$ the diagonal matrix $\operatorname{diag}\left(\lambda_{1}, \cdots, \lambda_{n}\right)$. We also denote by $\left\{a^{i j}\right\}$ and $\left\{c^{i j}\right\}$ the inverse matrices of $\left\{a_{i j}\right\}$ and $\left\{c_{i j}\right\}$, respectively.

We observe $g_{i j}(p)=\left\langle F_{i}, F_{j}\right\rangle(p)=\left\langle a_{i k} E_{k}, a_{j l} E_{l}\right\rangle=a_{i k} a_{j k}$. Also, we can obtain $F^{i}(p)=$ $a^{j i} E_{j}$ by $a^{j i}=a^{j k}\left\langle F_{k}(p), F^{i}(p)\right\rangle=\left\langle a^{j k} a_{k l} E_{l}, F^{i}(p)\right\rangle=\left\langle E_{j}, F^{i}(p)\right\rangle$. So, we have $g^{i j}(p)=$ $\left\langle F^{i}, F^{j}\right\rangle(p)=a^{k i} a^{k j}$. In addition, $L F_{i}(p)=h_{i j}(p) F^{j}(p)$ implies

$$
\begin{aligned}
a^{m i} E_{m} & =F^{i}(p)=b^{i j}(p) h_{j k}(p) F^{k}(p)=b^{i j}(p) L F_{j}(p) \\
& =b^{i j}(p) L\left(a_{j k} E_{k}\right)=b^{i j}(p) a_{j k} L\left(E_{k}\right)=b^{i j}(p) a_{j k} \lambda_{k} E_{k}=b^{i j}(p) a_{j k} c_{k m} E_{m}
\end{aligned}
$$

Hence, we have $b^{i n}(p)=b^{i j}(p) a_{j k} c_{k m} c^{m l} a^{l n}=a^{m i} c^{m l} a^{l n}$, and thus the following holds

$$
\begin{aligned}
b^{1 r} g_{r s} b^{s 1}(p) & =a^{i 1} c^{i j} a^{j r} a_{r k} a_{s k} a^{m 1} c^{m l} a^{l s}=a^{i 1} c^{i j} \delta_{k}^{j} \delta_{k}^{l} a^{m 1} c^{m l}=a^{i 1} c^{i j} a^{m 1} c^{m j}=\sum_{j}\left(a^{j 1}\right)^{2} \lambda_{j}^{-2} \\
& \leqslant \sum_{j}\left(a^{j 1}\right)^{2} \lambda_{\min }^{-2}=\lambda_{\min }^{-2} \sum_{k, j}\left\langle a^{k 1} E_{k}, a^{j 1} E_{j}\right\rangle=\lambda_{\min }^{-2}\left\langle F^{1}(p), F^{1}(p)\right\rangle=\lambda_{\min }^{-2} g^{11}(p),
\end{aligned}
$$

which is the desired result for $i=1$ and we can obtain the same result for each $i \in\{1, \cdots, n\}$.

## Chapter 3

## Regularity of strictly convex solutions

In this chapter, we study the regularity of strictly convex solutions. In particular, we will establish the lower bounds for principal curvatures and the upper bounds for the Gauss curvature, which yield the upper bounds for principal curvatures. By using the curvature estimates, the existence results [10], [12] were obtained as follows:

Theorem 3.1 (Chow). Let $\Sigma_{0}$ be a closed, smooth, and strictly convex hypersurface embedded in $\mathbb{R}^{n+1}$. Then, given an immersion $F_{0}: M^{n} \rightarrow \mathbb{R}^{n+1}$ such that $\Sigma_{0}=F_{0}\left(M^{n}\right)$, and for any $\alpha \in(0,+\infty)$, there exists a closed smooth and strictly convex solution $\Sigma_{t}:=F\left(M^{n}, t\right)$ of the $\alpha-$ Gauss curvature flow $\left.\sqrt{*^{\alpha}}\right)$ for finite time $0<t<T$, where $\Sigma_{t}$ converges to a point as $t$ approaches to $T$.

Theorem 3.2 (Choi-Daskalopoulos-Kim-Lee). Let $\Sigma_{0}$ be a complete, non-compact, and strictly convex hypersurface embedded in $\mathbb{R}^{n+1}$. Then, given an immersion $F_{0}: M^{n} \rightarrow \mathbb{R}^{n+1}$ such that $\Sigma_{0}=F_{0}\left(M^{n}\right)$, and for any $\alpha \in(0,+\infty)$, there exists a complete, non-compact, smooth and strictly convex solution $\Sigma_{t}:=F\left(M^{n}, t\right)$ of the $\alpha$-Gauss curvature flow $\left(*^{\alpha}\right)$ which is defined for all time $0<t<+\infty$.


Figure 1. Examples of the initial graph $\Sigma_{0}$ defined on a domain $\Omega \subset \mathbb{R}^{n}$.

## 1 Lower bounds on the principal curvatures

In this section we will obtain lower bounds on the principal curvatures. One can easily establish the following lower bounds by combining equations (2.1) and 2.8) and applying the maximum principle.

Theorem 3.3. Given a closed, smooth, and strictly convex solution $\Sigma_{t}$ to $\left({ }^{*}\right)$, the following holds

$$
b^{i j} g_{i j} \leqslant \max _{\Sigma_{0}} b^{i j} g_{i j}
$$

We will achieve the interior estimates by using Pogorelov type estimates with respect to $b^{i i}$. Recall that $b^{i j}$ denotes the inverse matrix of the second fundamental form $h_{i j}$. Since $b^{i i}$ depends on charts, we will find the relation between $b^{i i}$ and the principal curvatures as $[\mathbf{8}]$. We recall the definition of the cut-off function $\psi_{\beta}(p, t):=(M-\beta t-\bar{u}(p, t))_{+}$, where $\bar{u}(p, t):=\left\langle F(p, t), \vec{e}_{n+1}\right\rangle$ denotes the height function.

Theorem 3.4 (Local lower bound for the principal curvatures). Assume that $\Sigma_{t}$ is a complete strictly convex smooth graph solution of $(\sqrt[*^{\alpha}]{ })$ defined on $M^{n} \times[0, T]$, for some $T>0$. Then, given constants $\beta>0$ and $M \geqslant \beta$, the following holds

$$
\left(\psi_{\beta}^{-n\left(1+\frac{1}{\alpha}\right)} \lambda_{\min }\right)(p, t) \geqslant M^{-n\left(1+\frac{1}{\alpha}\right)} \min \left\{\inf _{Q_{M}} \lambda_{\min }(p, 0), \frac{\beta^{(n-1)+\frac{1}{\alpha}}}{\left(n^{2}(n+1)(n \alpha-1)_{+}\right)^{(n-1)+\frac{1}{\alpha}}}\right\}
$$

where $Q_{M}=\left\{p \in M^{n}: \bar{u}(p, 0)<M\right\}$, and $(n \alpha-1)_{+}^{-1}=+\infty$, if $n \alpha \leqslant 1$.
Proof. Consider the cut-off function $\psi_{\beta}:=(M-\beta t-\bar{u}(p, t))_{+}$, where $\bar{u}(p, t)=\left\langle F(p, t), \vec{e}_{n+1}\right\rangle$ denotes the height function. Since $\Sigma_{t}$ is a complete and strictly convex graph, $\psi_{\beta}$ is compactly supported. Therefore, for a fixed $T \in(0,+\infty)$, the function $\psi_{\beta}^{n\left(1+\frac{1}{\alpha}\right)} \lambda_{\min }^{-1}$ attains its maximum on $M^{n} \times[0, T]$ at a point $\left(p_{0}, t_{0}\right)$. If $t_{0}=0$, then we obtain the desired result by the bound $\psi_{\beta} \leqslant M$ and the conditions on our initial data. So, we may assume in what follows that $t_{0}>0$.

We begin by choosing a chart $(U, \varphi)$ with $p_{0} \in \varphi(U) \subset M^{n}$, on which the covariant derivatives $\left\{\nabla_{i} F\left(p_{0}, t_{0}\right):=\partial_{i}(F \circ \varphi)\left(\varphi^{-1}\left(p_{0}\right), t_{0}\right)\right\}_{i=1, \cdots, n}$ form an orthonormal basis of $\left(T \Sigma_{t_{0}}\right)_{F\left(p_{0}, t_{0}\right)}$ satisfying

$$
g_{i j}\left(p_{0}, t_{0}\right)=\delta_{i j}, \quad h_{i j}\left(p_{0}, t_{0}\right)=\delta_{i j} \lambda_{i}\left(p_{0}, t_{0}\right), \quad \lambda_{1}\left(p_{0}, t_{0}\right)=\lambda_{\min }\left(p_{0}, t_{0}\right) .
$$

In particular, at the point $\left(p_{0}, t_{0}\right)$ we have $b^{11}\left(p_{0}, t_{0}\right)=\lambda_{\min }^{-1}\left(p_{0}, t_{0}\right)$ and $g^{11}\left(p_{0}, t_{0}\right)=1$. Next, we define the function $w: \varphi(U) \times[0, T] \rightarrow \mathbb{R}$ by

$$
w:=\psi_{\beta}^{n\left(1+\frac{1}{\alpha}\right)} \frac{b^{11}}{g^{11}} .
$$

Notice that on the chart $(U, \varphi)$, if $t \neq t_{0}$, then the covariant derivatives $\left\{\nabla_{i} F\left(p_{0}, t\right)\right\}_{i=1, \cdots, n}$ may not form an orthonormal basis of $\left(T \Sigma_{t}\right)_{F\left(p_{0}, t\right)}$. However, since Proposition 2.4 holds for every chart and immersion, we have $w \leqslant \psi_{\beta}^{n\left(1+\frac{1}{\alpha}\right)} \lambda_{\text {min }}^{-1}$. Hence, for $(p, t) \in \varphi(U) \times[0, T]$, the following holds

$$
w(p, t) \leqslant \psi_{\beta}^{n\left(1+\frac{1}{\alpha}\right)} \lambda_{\min }^{-1}(p, t) \leqslant \psi_{\beta}^{n\left(1+\frac{1}{\alpha}\right)} \lambda_{\min }^{-1}\left(p_{0}, t_{0}\right)=w\left(p_{0}, t_{0}\right)
$$

which shows that $w$ attains its maximum at $\left(p_{0}, t_{0}\right)$.
Observe next that since $\nabla g^{11}=0$, the following holds on the support of $\psi_{\beta}$

$$
\begin{equation*}
\frac{\nabla_{i} w}{w}=n\left(1+\frac{1}{\alpha}\right) \frac{\nabla_{i} \psi_{\beta}}{\psi_{\beta}}+\frac{\nabla_{i} b^{11}}{b^{11}} . \tag{3.1}
\end{equation*}
$$

Let us differentiate the equation above, again.

$$
\frac{\nabla_{i} \nabla_{j} w}{w}-\frac{\nabla_{i} w \nabla_{j} w}{w^{2}}=n\left(1+\frac{1}{\alpha}\right) \frac{\nabla_{i} \nabla_{j} \psi_{\beta}}{\psi_{\beta}}-n\left(1+\frac{1}{\alpha}\right) \frac{\nabla_{i} \psi_{\beta} \nabla_{j} \psi_{\beta}}{\psi_{\beta}^{2}}+\frac{\nabla_{i} \nabla_{j} b^{11}}{b^{11}}-\frac{\nabla_{i} b^{11} \nabla_{j} b^{11}}{\left(b^{11}\right)^{2}} .
$$

Multiply by $\alpha K^{\alpha} b^{i j}$ and sum the equations over all $i, j$ to obtain

$$
\frac{\mathcal{L} w}{w}-\frac{\|\nabla w\|_{\mathcal{L}}^{2}}{w^{2}}=n\left(1+\frac{1}{\alpha}\right) \frac{\mathcal{L} \psi_{\beta}}{\psi_{\beta}}-n\left(1+\frac{1}{\alpha}\right) \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}+\frac{\mathcal{L} b^{11}}{b^{11}}-\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}
$$

On the other hand, on the support of $\psi_{\beta}$, we also have

$$
\frac{\partial_{t} w}{w}=n\left(1+\frac{1}{\alpha}\right) \frac{\partial_{t} \psi_{\beta}}{\psi_{\beta}}+\frac{\partial_{t} b^{11}}{b^{11}}-\frac{\partial_{t} g^{11}}{g^{11}} .
$$

Subtract the equations above. Then, $w^{-2}\|\nabla w\|_{\mathcal{L}}^{2} \geqslant 0$ implies the following inequality

$$
\frac{\mathcal{L} w}{w}-\frac{\partial_{t} w}{w} \geqslant n\left(1+\frac{1}{\alpha}\right)\left(\frac{\mathcal{L} \psi_{\beta}}{\psi_{\beta}}-\frac{\partial_{t} \psi_{\beta}}{\psi_{\beta}}\right)-n\left(1+\frac{1}{\alpha}\right) \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}+\left(\frac{\mathcal{L} b^{11}}{b^{11}}-\frac{\partial_{t} b^{11}}{b^{11}}\right)-\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}+\frac{\partial_{t} g^{11}}{g^{11}}
$$

By 2.2 and (2.4) we have $\partial_{t} g^{11}=2 K^{\alpha} h^{11}$ and $\mathcal{L} \psi_{\beta}-\partial_{t} \psi_{\beta}=\beta-(n \alpha-1) v^{-1} K^{\alpha}$, while by (2.8)

$$
\mathcal{L} b^{11}-\partial_{t} b^{11}=\alpha K^{\alpha} b^{i 1} b^{j 1}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right) \nabla_{i} h_{k l} \nabla_{j} h_{m n}+\alpha K^{\alpha} H b^{11}-(1+n \alpha) K^{\alpha} g^{11} .
$$

Combining the above yields

$$
\begin{align*}
& \frac{\mathcal{L} w}{w}-\frac{\partial_{t} w}{w} \geqslant-n\left(1+\frac{1}{\alpha}\right) \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}-\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}+\frac{\alpha K^{\alpha} b^{i 1} b^{j 1}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right)}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} \\
&  \tag{3.2}\\
& \quad-n\left(1+\frac{1}{\alpha}\right)(n \alpha-1) \frac{K^{\alpha} v^{-1}}{\psi_{\beta}}+n\left(1+\frac{1}{\alpha}\right) \frac{\beta}{\psi_{\beta}}+\alpha K^{\alpha} H-(1+n \alpha) K^{\alpha} \frac{g^{11}}{b^{11}}+2 K^{\alpha} \frac{h^{11}}{g^{11}} .
\end{align*}
$$

Now, at $\left(p_{0}, t_{0}\right)$, the following holds

$$
\begin{equation*}
\alpha K^{\alpha} H-(1+n \alpha) K^{\alpha} \frac{g^{11}}{b^{11}}+2 K^{\alpha} \frac{h^{11}}{g^{11}} \geqslant n \alpha K^{\alpha} \lambda_{\min }-(1+n \alpha) K^{\alpha} \lambda_{\min }+2 K^{\alpha} \lambda_{\min }=K^{\alpha} \lambda_{\min } . \tag{3.3}
\end{equation*}
$$

In addition, if $n \alpha \geqslant 1$, then $H \geqslant n \lambda_{\text {min }}$ gives

$$
\alpha K^{\alpha} H=\left(\alpha-\frac{1}{n}\right) K^{\alpha} H+\frac{1}{n} K^{\alpha} H \geqslant(n \alpha-1) K^{\alpha} \lambda_{\min }+\frac{1}{n} K^{\alpha} H .
$$

Therefore, in the case that $n \alpha \geqslant 1$, we can improve (3.3) to obtain

$$
\begin{equation*}
\alpha K^{\alpha} H-(1+n \alpha) K^{\alpha} \frac{g^{11}}{b^{11}}+2 K^{\alpha} \frac{h^{11}}{g^{11}} \geqslant \frac{1}{n} K^{\alpha} H . \tag{3.4}
\end{equation*}
$$

Also, at the maximum point $\left(p_{0}, t_{0}\right)$ of $w, \nabla w\left(p_{0}, t_{0}\right)=0$ holds. So, 3.1 leads to $\frac{n(1+\alpha)}{\alpha} \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}+\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}=\left(1+\frac{\alpha}{n(1+\alpha)}\right) \frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}=\left(1+\frac{\alpha}{n(1+\alpha)}\right) \alpha \sum_{i=1}^{n} \frac{b^{i i} K^{\alpha}\left|\nabla_{i} b^{11}\right|^{2}}{\left(b^{11}\right)^{2}}$.
From (2.13), we get $\nabla_{i} b^{11}=-b^{1 k} b^{1 /} \nabla_{i} h_{k l}=-\left(b^{11}\right)^{2} \nabla_{i} h_{11}$ at $\left(p_{0}, t_{0}\right)$. Hence,

$$
\frac{n(1+\alpha)}{\alpha} \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}+\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}=\alpha\left(1+\frac{\alpha}{n(1+\alpha)}\right) \sum_{i=1}^{n} b^{i i}\left(b^{11}\right)^{2} K^{\alpha}\left|\nabla_{i} h_{11}\right|^{2}
$$

Defining, at $\left(p_{0}, t_{0}\right)$, we define

$$
\begin{aligned}
& I_{i}=b^{i i}\left(b^{11}\right)^{2} K^{\alpha}\left|\nabla_{i} h_{11}\right|^{2} \\
& J_{i}=b^{i i} \nabla_{1} h_{i i} .
\end{aligned}
$$

We may rewrite the equation above as

$$
\begin{equation*}
\frac{n(1+\alpha)}{\alpha} \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}+\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}=\alpha\left(1+\frac{\alpha}{n(1+\alpha)}\right) \sum_{i=1}^{n} I_{i} \tag{3.5}
\end{equation*}
$$

and also write

$$
b^{k l} b^{m n} \nabla_{1} h_{k l} \nabla_{1} h_{m n}=\left|b^{m n} \nabla_{1} h_{m n}\right|^{2}=\left|\sum_{i=1}^{n} b^{i i} \nabla_{1} h_{i i}\right|^{2}=\left|\sum_{i=1}^{n} J_{i}\right|^{2}
$$

which gives

$$
\frac{\alpha^{2} K^{\alpha} b^{i 1} b^{j 1} b^{k l} b^{m n}}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n}=\alpha^{2} K^{\alpha} b^{11} b^{k l} b^{m n} \nabla_{1} h_{k l} \nabla_{1} h_{m n}=\alpha^{2} K^{\alpha} b^{11}\left|\sum_{i=1}^{n} J_{i}\right|^{2} .
$$

Since $\alpha^{2} \geqslant \frac{\alpha^{2}}{1+\alpha}$, we conclude that

$$
\begin{equation*}
\frac{\alpha^{2} K^{\alpha} b^{i 1} b^{j 1} b^{k l} b^{m n}}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} \geqslant \frac{\alpha^{2}}{1+\alpha} K^{\alpha} b^{11}\left|\sum_{i=1}^{n} J_{i}\right|^{2} . \tag{3.6}
\end{equation*}
$$

Finally, at ( $p_{0}, t_{0}$ ), we also have

$$
\begin{aligned}
\frac{\alpha K^{\alpha} b^{i 1} b^{j 1} b^{k m} b^{n l}}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} & =\alpha K^{\alpha} b^{11} b^{k m} b^{n l} \nabla_{1} h_{k l} \nabla_{1} h_{m n} \\
& =\alpha K^{\alpha} b^{11}\left(\sum_{i=1}^{n}\left|b^{i i} \nabla_{1} h_{i i}\right|^{2}+\sum_{i \neq j} b^{i i} b^{j j}\left|\nabla_{1} h_{i j}\right|^{2}\right) \\
& \geqslant \alpha K^{\alpha} b^{11}\left(\sum_{i=1}^{n}\left|b^{i i} \nabla_{1} h_{i i}\right|^{2}+2 \sum_{i \neq 1} b^{i i} b^{11}\left|\nabla_{1} h_{i 1}\right|^{2}\right) \\
& =\alpha K^{\alpha} b^{11} \sum_{i=1}^{n}\left|J_{i}\right|^{2}+2 \alpha \sum_{i \neq 1} I_{i} .
\end{aligned}
$$

Using $\alpha \geqslant \frac{\alpha^{2}}{1+\alpha}$, we may rewrite the inequality above as

$$
\begin{equation*}
\frac{\alpha K^{\alpha} b^{i 1} b^{j 1} b^{k m} b^{n l}}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} \geqslant \frac{\alpha^{2}}{1+\alpha} K^{\alpha} b^{11} \sum_{i \neq 1}\left|J_{i}\right|^{2}+\alpha K^{\alpha} b^{11}\left|J_{1}\right|^{2}+2 \alpha \sum_{i \neq 1} I_{i} . \tag{3.7}
\end{equation*}
$$

Adding (3.6) and (3.7), gives that at $\left(p_{0}, t_{0}\right)$ we have

$$
\frac{\alpha K^{\alpha} b^{i 1} b^{j 1}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right)}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} \geqslant \frac{\alpha^{2}}{1+\alpha} K^{\alpha} b^{11}\left(\left|\sum_{i=1}^{n} J_{i}\right|^{2}+\sum_{i \neq 1}\left|J_{i}\right|^{2}\right)+\alpha K^{\alpha} b^{11}\left|J_{1}\right|^{2}+2 \alpha \sum_{i \neq 1} I_{i}
$$

and by the Cauchy-Schwarz inequality

$$
\begin{aligned}
n\left(\left|\sum_{i=1}^{n} J_{i}\right|^{2}+\sum_{i \neq 1}\left|J_{i}\right|^{2}\right) & =\left(1^{2}+(-1)^{2}+\cdots+(-1)^{2}\right)\left(\left|\sum_{i=1}^{n} J_{i}\right|^{2}+\sum_{i \neq 1}\left|J_{i}\right|^{2}\right) \\
& \geqslant\left|\sum_{i=1}^{n} J_{i}+\sum_{i \neq 1}-J_{i}\right|^{2}=\left|J_{1}\right|^{2}
\end{aligned}
$$

and $2 \alpha \geqslant \alpha\left(1+\frac{\alpha}{n(1+\alpha)}\right)$, we obtain

$$
\frac{\alpha K^{\alpha} b^{i 1} b^{j 1}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right)}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} \geqslant \alpha\left(1+\frac{\alpha}{n(1+\alpha)}\right)\left(K^{\alpha} b^{11}\left|J_{1}\right|^{2}+\sum_{i \neq 1} I_{i}\right) .
$$

Combining the last inequality with (3.5) while using that $K^{\alpha} b^{11}\left|J_{1}\right|^{2}=I_{1}$, yields

$$
\begin{equation*}
-\frac{n(1+\alpha)}{\alpha} \frac{\left\|\nabla \psi_{\beta}\right\|_{\mathcal{L}}^{2}}{\psi_{\beta}^{2}}-\frac{\left\|\nabla b^{11}\right\|_{\mathcal{L}}^{2}}{\left(b^{11}\right)^{2}}+\frac{\alpha K^{\alpha} b^{i 1} b^{j 1}\left(\alpha b^{k l} b^{m n}+b^{k m} b^{n l}\right)}{b^{11}} \nabla_{i} h_{k l} \nabla_{j} h_{m n} \geqslant 0 . \tag{3.8}
\end{equation*}
$$

We conclude by (3.2), (3.3) and (3.8) that at $\left(p_{0}, t_{0}\right)$ the following holds

$$
0 \geqslant \frac{\mathcal{L} w}{w}-\frac{\partial_{t} w}{w} \geqslant-n\left(1+\frac{1}{\alpha}\right)(n \alpha-1) \frac{K^{\alpha} v^{-1}}{\psi_{\beta}}+n\left(1+\frac{1}{\alpha}\right) \frac{\beta}{\psi_{\beta}}+K^{\alpha} \lambda_{\text {min }}
$$

If $n \alpha \leqslant 1$, then the last inequality gives $n\left(1+\frac{1}{\alpha}\right) \leqslant 0$, a contradiction. Hence $t_{0}=0$, and therefore the desired result holds. If $n \alpha \geqslant 1$, then we use the improved inequality (3.4) instead of (3.3) by and perform the same estimates for all the other terms, so that at $\left(p_{0}, t_{0}\right)$, we obtain

$$
0 \geqslant \frac{\mathcal{L} w}{w}-\frac{\partial_{t} w}{w} \geqslant-n\left(1+\frac{1}{\alpha}\right)(n \alpha-1) \frac{K^{\alpha} v^{-1}}{\psi_{\beta}}+n\left(1+\frac{1}{\alpha}\right) \frac{\beta}{\psi_{\beta}}+\frac{1}{n} K^{\alpha} H
$$

Hence

$$
n\left(1+\frac{1}{\alpha}\right)(n \alpha-1) \frac{v^{-1}}{\psi_{\beta}} \geqslant n\left(1+\frac{1}{\alpha}\right) \frac{\beta}{\psi_{\beta}} K^{-\alpha}+\frac{1}{n} H .
$$

Since $n \alpha \geqslant 1$, we have $1+\frac{1}{\alpha} \leqslant 1+n$, and using also that $v \geqslant 1, \psi_{\beta} \leqslant M$, and $M \geqslant \beta$, we conclude from the previous inequality that

$$
n^{2}(1+n)(n \alpha-1) \psi_{\beta}^{-1} \geqslant n^{2}\left(1+\frac{1}{\alpha}\right) K^{-\alpha} \frac{\beta}{\psi_{\beta}}+H \geqslant\left(n^{2}\left(1+\frac{1}{\alpha}\right) K^{-\alpha}+H\right) \frac{\beta}{M}
$$

Next, we employ the Young's inequality

$$
\frac{K^{-\alpha}}{(n-1) \alpha+1}+\frac{(n-1) \alpha H}{(n-1) \alpha+1} \geqslant K^{-\frac{\alpha}{(n-1) \alpha+1}} H^{\frac{(n-1) \alpha}{(n-1) \alpha+1}}=\left(K^{-1} H^{n-1}\right)^{\frac{\alpha}{(n-1) \alpha+1}}
$$

and observe the following

$$
K^{-1} H^{n-1}=\frac{H^{n-1}}{\lambda_{1} \lambda_{2} \cdots \lambda_{n}}=\frac{1}{\lambda_{1}} \frac{H^{n-1}}{\lambda_{2} \cdots \lambda_{n}} \geqslant \frac{1}{\lambda_{1}}=\lambda_{\min }^{-1} .
$$

Combining the last three inequalities yields
$n^{2}(1+n)(n \alpha-1) M \beta^{-1} \psi_{\beta}^{-1} \geqslant n^{2}\left(1+\frac{1}{\alpha}\right) K^{-\alpha}+H \geqslant \frac{K^{-\alpha}}{(n-1) \alpha+1}+\frac{(n-1) \alpha H}{(n-1) \alpha+1} \geqslant\left(\lambda_{\min }^{-1}\right)^{\frac{\alpha}{n-1) \alpha+1}}$.
We conclude, that if $n \alpha \geqslant 1$, the following holds at ( $p_{0}, t_{0}$ )

$$
\lambda_{\min }^{-1} \psi_{\beta}^{n\left(1+\frac{1}{\alpha}\right)} \leqslant\left(n^{2}(n+1)(n \alpha-1) M \beta^{-1}\right)^{(n-1)+\frac{1}{\alpha}} \psi_{\beta}^{1+\frac{n-1}{\alpha}}
$$

Thus, $\psi_{\beta} \leqslant M$ gives the desired result.

## 2 Speed estimates

We establish upper bounds for the speed $K^{\alpha}$. Andrews obtained the following result for closed solutions to applying the maximum principle to the quantity $t^{\frac{n \alpha}{\alpha \alpha+1}} K^{\alpha}(2\langle F, \vec{n}\rangle-\rho)^{-1}$ in [2].

Theorem 3.5 (Andrews). Let $\Sigma_{t}$ be a closed, smooth and strictly convex solution to $\left(*^{\alpha}\right)$ defined for $t \in(0, T]$. Assume that $\Sigma_{T}$ encloses a $(n+1)$-ball $B_{\rho}^{n+1}(0)$ of radius $\rho>0$ and $\Sigma_{0}$ is enclosed by $B_{R}^{n+1}(0)$. Then, the following holds

$$
t^{\frac{n \alpha}{n \alpha+1}} K^{\alpha} \leqslant C(R, \rho, T, \alpha, n),
$$

for some constant $C(R, \rho, T, \alpha, n)$ depending on $R, \rho, T, \alpha, n$.

To establish the interior estimate, we first derive the gradient estimate. We recall the notation $v:=\left\langle\vec{n}, \vec{e}_{n+1}\right\rangle^{-1}$ (gradient function) and $\psi_{\beta}(p, t):=(M-\beta t-\bar{u}(p, t))_{+}$(cut-off function) where $\bar{u}(p, t)=\left\langle F(p, t), \vec{e}_{n+1}\right\rangle$ denotes the height function.

Theorem 3.6 (Gradient estimate). Assume that $\Sigma_{t}$ is a complete strictly convex smooth graph solution of $(\sqrt[*^{\alpha}]{ })$ defined on $M^{n} \times[0, T]$, for some $T>0$. Given constants $\beta>0$ and $M \geqslant \beta$,

$$
v(p, t) \psi_{\beta}(p, t) \leqslant M \max \left\{\sup _{Q_{M}} v(p, 0), \beta^{-1} n^{\frac{1}{n \alpha+1}}(n \alpha-1)_{+}\right\}
$$

where $Q_{M}=\left\{p \in M^{n}: \bar{u}(p, 0)<M\right\}$.
Proof. First use (2.4) and (2.9), that is

$$
\begin{aligned}
\partial_{t} \psi_{\beta} & =\mathcal{L} \psi_{\beta}+(n \alpha-1) v^{-1} K^{\alpha}-\beta \\
\partial_{t} v & =\mathcal{L} v-2 v^{-1}\|\nabla v\|_{\mathcal{L}}^{2}-\alpha K^{\alpha} H v
\end{aligned}
$$

to compute

$$
\begin{aligned}
\partial_{t}\left(\psi_{\beta} v\right) & =\psi_{\beta} \mathcal{L} v-2 \psi_{\beta} v^{-1}\|\nabla v\|_{\mathcal{L}}^{2}-\alpha \psi_{\beta} K^{\alpha} H v+v \mathcal{L} \psi_{\beta}+(n \alpha-1) K^{\alpha}-\beta v \\
& =\mathcal{L}\left(\psi_{\beta} v\right)-2\left\langle\nabla \psi_{\beta}, \nabla v\right\rangle_{\mathcal{L}}-2 \psi_{\beta} v^{-1}\|\nabla v\|_{\mathcal{L}}^{2}-\alpha \psi_{\beta} K^{\alpha} H v+(n \alpha-1) K^{\alpha}-\beta v \\
& =\mathcal{L}\left(\psi_{\beta} v\right)-2 v^{-1}\left\langle\nabla\left(\psi_{\beta} v\right), \nabla v\right\rangle_{\mathcal{L}}-\alpha \psi_{\beta} K^{\alpha} H v+(n \alpha-1) K^{\alpha}-\beta v .
\end{aligned}
$$

Since $\psi_{\beta}$ is compactly supported, for a fixed $T \in(0,+\infty)$ the function $\psi_{\beta} v$ attains its maximum on $M^{n} \times[0, T]$ at some $\left(p_{0}, t_{0}\right)$. If $t_{0}=0$, then we obtain the desired result. Assume that $t_{0}>0$. Then, at $\left(p_{0}, t_{0}\right)$, we have

$$
\partial_{t}\left(\psi_{\beta} v\right)-\mathcal{L}\left(\psi_{\beta} v\right) \geqslant 0
$$

which is equivalent to

$$
(n \alpha-1) K^{\alpha} \geqslant \alpha \psi_{\beta} K^{\alpha} H v+\beta v=\left(\alpha \psi_{\beta} K^{\alpha} H+\beta\right) v
$$

Hence, an interior maximum can be achieved only if $n \alpha \geqslant 1$. From now on, we assume that $n \alpha \geqslant 1$. If we multiply the last inequality by $M K^{-\alpha}$ and use that $M \geqslant \beta$ we obtain

$$
(n \alpha-1) M \geqslant\left(\alpha M \psi_{\beta} H+\beta M K^{-\alpha}\right) v \geqslant \beta\left(\alpha \psi_{\beta} H+M K^{-\alpha}\right) v .
$$

On the other hand, $M \geqslant \psi_{\beta}$ implies $\alpha \psi_{\beta} H+M K^{-\alpha} \geqslant \alpha \psi_{\beta} H+\psi_{\beta} K^{-\alpha} \geqslant\left(\alpha H+K^{-\alpha}\right) \psi_{\beta}$. Hence,

$$
(n \alpha-1) M \geqslant \beta\left(\alpha H+K^{-\alpha}\right) \psi_{\beta} v \geqslant \frac{\beta}{n}\left(\frac{n \alpha}{n \alpha+1} H+\frac{1}{n \alpha+1} K^{-\alpha}\right) \psi_{\beta} v .
$$

Next, we apply the Young's inequality

$$
\frac{n \alpha}{n \alpha+1} H+\frac{1}{n \alpha+1} K^{-\alpha} \geqslant H^{\frac{n \alpha}{n \alpha+1}} K^{-\frac{\alpha}{n \alpha+1}}=\left(H K^{-\frac{1}{n}}\right)^{\frac{n \alpha}{n \alpha+1}} .
$$

Since $H \geqslant n K^{1 / n}$, we conclude that in the case $n \alpha \geqslant 1, t_{0}>0$, the following holds at ( $p_{0}, t_{0}$ )

$$
(n \alpha-1)_{+} M \geqslant \beta n^{-1} n^{\frac{n \alpha}{n \alpha+1}} \psi_{\beta} v=\beta n^{-\frac{1}{n \alpha+1}} \psi_{\beta} v
$$

from which the desired inequality readily follows.

We now apply the trick developed in [6].
Theorem 3.7 (Local speed bound). Assume that $\Sigma_{t}$ is a complete strictly convex smooth graph solution of $\left(*^{\alpha}\right)$ defined on $M^{n} \times[0, T)$. Then, given a constant $M \geqslant 1$,

$$
\left(\frac{t}{1+t}\right)\left(\psi^{2} K^{\frac{1}{n}}\right)(p, t) \leqslant(4 n \alpha+1)^{2}(2 \theta)^{1+\frac{1}{2 n \alpha}}\left(\theta \Lambda+M^{2}\right)
$$

where $\theta$ and $\Lambda$ are constants given by

$$
\begin{aligned}
\theta & =\sup \left\{v^{2}(p, s): \bar{u}(p, s)<M, s \in[0, t]\right\} \\
\Lambda & =\sup \left\{\lambda_{\min }^{-1}(p, s): \bar{u}(p, s)<M, s \in[0, t]\right\}
\end{aligned}
$$

Proof. Choosing a fixed time $T_{0} \in(0, T)$, we redefine $\theta$ and $\Lambda$ by

$$
\theta=\sup \left\{v^{2}(p, t): \bar{u}(p, t)<M, t \in\left[0, T_{0}\right]\right\}, \quad \Lambda=\sup \left\{\lambda_{\min }^{-1}(p, t): \bar{u}(p, t)<M, t \in\left[0, T_{0}\right]\right\}
$$

Also, we define $\eta:[0, T) \rightarrow \mathbb{R}$ by

$$
\eta(t)=\frac{t}{1+t}
$$

which will be used later in this proof.
Following the well used idea by Caffarelli, Nirenberg and Spruck in [6] (see also in [17] and [8]), we define the function $\varphi=\varphi\left(v^{2}\right)$, depending on $v^{2}$, by

$$
\varphi\left(v^{2}\right)=\frac{v^{2}}{2 \theta-v^{2}}
$$

The evolution equation of $v$ in (2.9) gives

$$
\partial_{t}\left(v^{2}\right)=\mathcal{L}\left(v^{2}\right)-2 \alpha K^{\alpha} H v^{2}-6\|\nabla v\|_{\mathcal{L}}^{2} .
$$

Then, the evolution equation of $\varphi\left(v^{2}\right)$ is

$$
\partial_{t} \varphi=\varphi^{\prime}\left(\mathcal{L} v^{2}-2 \alpha K^{\alpha} H v^{2}-6\|\nabla v\|_{\mathcal{L}}^{2}\right)=\mathcal{L} \varphi-\varphi^{\prime \prime}\left\|\nabla v^{2}\right\|_{\mathcal{L}}^{2}-\varphi^{\prime}\left(2 \alpha K^{\alpha} H v^{2}+6\|\nabla v\|_{\mathcal{L}}^{2}\right) .
$$

Also, the evolution equation of $K^{\alpha}$ in (2.7) leads to

$$
\partial_{t} K^{2 \alpha}=\mathcal{L} K^{2 \alpha}-\frac{1}{2} K^{-2 \alpha}\left\|\nabla K^{2 \alpha}\right\|_{\mathcal{L}}^{2}+2 \alpha K^{3 \alpha} H
$$

implying the following evolution equation for $K^{2 \alpha} \varphi\left(v^{2}\right)$

$$
\begin{aligned}
\partial_{t}\left(K^{2 \alpha} \varphi\right)= & \mathcal{L}\left(K^{2 \alpha} \varphi\right)-2\left\langle\nabla K^{2 \alpha}, \nabla \varphi\right\rangle_{\mathcal{L}}+2 \alpha K^{3 \alpha} H\left(\varphi-\varphi^{\prime} v^{2}\right) \\
& -\frac{1}{2} \varphi K^{-2 \alpha}\left\|\nabla K^{2 \alpha}\right\|_{\mathcal{L}}^{2}-\left(4 \varphi^{\prime \prime} v^{2}+6 \varphi^{\prime}\right) K^{2 \alpha}\|\nabla v\|_{\mathcal{L}}^{2} .
\end{aligned}
$$

Observe that

$$
\begin{aligned}
-2\left\langle\nabla K^{2 \alpha}, \nabla \varphi\right\rangle_{\mathcal{L}} & =-\left\langle\nabla K^{2 \alpha}, \nabla \varphi\right\rangle_{\mathcal{L}}+\varphi^{-1} K^{2 \alpha}\|\nabla \varphi\|_{\mathcal{L}}^{2}-\varphi^{-1}\left\langle\nabla \varphi, \nabla\left(K^{2 \alpha} \varphi\right)\right\rangle_{\mathcal{L}} \\
& \leqslant \frac{1}{2} \varphi K^{-2 \alpha}\left\|\nabla K^{2 \alpha}\right\|_{\mathcal{L}}^{2}+\frac{3}{2} \varphi^{-1} K^{2 \alpha}\|\nabla \varphi\|_{\mathcal{L}}^{2}-\varphi^{-1}\left\langle\nabla \varphi, \nabla\left(K^{2 \alpha} \varphi\right)\right\rangle_{\mathcal{L}} .
\end{aligned}
$$

Hence, the following inequality holds

$$
\begin{align*}
\partial_{t}\left(K^{2 \alpha} \varphi\right) \leqslant & \mathcal{L}\left(K^{2 \alpha} \varphi\right)-\varphi^{-1}\left\langle\nabla \varphi, \nabla\left(K^{2 \alpha} \varphi\right)\right\rangle_{\mathcal{L}}+2 \alpha K^{3 \alpha} H\left(\varphi-\varphi^{\prime} v^{2}\right)  \tag{3.9}\\
& -\left(4 \varphi^{\prime \prime} v^{2}+6 \varphi^{\prime}-6 \varphi^{-1} \varphi^{\prime 2} v^{2}\right) K^{2 \alpha}\|\nabla v\|_{\mathcal{L}}^{2}
\end{align*}
$$

Now, we have

$$
\varphi\left(v^{2}\right)+1=\frac{2 \theta}{2 \theta-v^{2}}, \quad \varphi^{\prime}\left(v^{2}\right)=\frac{2 \theta}{\left(2 \theta-v^{2}\right)^{2}}, \quad \varphi^{\prime \prime}\left(v^{2}\right)=\frac{4 \theta}{\left(2 \theta-v^{2}\right)^{3}} .
$$

Therefore, by direct calculation we obtain

$$
\varphi-\varphi^{\prime} v^{2}=\frac{v^{2}}{2 \theta-v^{2}}-\frac{2 \theta v^{2}}{\left(2 \theta-v^{2}\right)^{2}}=-\frac{v^{4}}{\left(2 \theta-v^{2}\right)^{2}}=-\varphi^{2}
$$

and

$$
\varphi^{-1} \nabla \varphi=\frac{2 \theta-v^{2}}{v^{2}} \frac{4 \theta v \nabla v}{\left(2 \theta-v^{2}\right)^{2}}=4 \theta \varphi v^{-3} \nabla v
$$

and also

$$
4 \varphi^{\prime \prime} v^{2}+6 \varphi^{\prime}-6 \varphi^{-1} \varphi^{\prime 2} v^{2}=\frac{16 \theta v^{2}}{\left(2 \theta-v^{2}\right)^{3}}+\frac{12 \theta}{\left(2 \theta-v^{2}\right)^{2}}-\frac{24 \theta^{2}}{\left(2 \theta-v^{2}\right)^{3}}=\frac{4 \theta}{\left(2 \theta-v^{2}\right)^{2}} \varphi
$$

Setting $f:=K^{2 \alpha} \varphi$ in 3.9) and using the identities above yields

$$
\partial_{t} f \leqslant \mathcal{L} f-4 \theta \varphi v^{-3}\langle\nabla v, \nabla f\rangle_{\mathcal{L}}-2 \alpha f K^{\alpha} H \varphi-\frac{4 \theta}{\left(2 \theta-v^{2}\right)^{2}} f\|\nabla v\|_{\mathcal{L}}^{2}
$$

On the other hand, (2.4) gives

$$
\partial_{t} \psi=\mathcal{L} \psi+(n \alpha-1) v^{-1} K^{\alpha} \leqslant \mathcal{L} \psi+n \alpha K^{\alpha} .
$$

Hence, on the support of $\psi$, we have

$$
\partial_{t} \psi^{4 n \alpha} \leqslant \mathcal{L} \psi^{4 n \alpha}-4 n \alpha(4 n \alpha-1) \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2}+4 n^{2} \alpha^{2} K^{\alpha} \psi^{4 n \alpha-1}
$$

Thus, on the support of $\psi$, the following holds

$$
\begin{aligned}
\partial_{t}\left(f \psi^{4 n \alpha}\right) \leqslant & \mathcal{L}\left(f \psi^{4 n \alpha}\right)-2\left\langle\nabla \psi^{4 n \alpha}, \nabla f\right\rangle_{\mathcal{L}}-4 \theta \varphi v^{-3} \psi^{4 n \alpha}\langle\nabla v, \nabla f\rangle_{\mathcal{L}}-2 \alpha f K^{\alpha} H \varphi \psi^{4 n \alpha} \\
& -\frac{4 \theta}{\left(2 \theta-v^{2}\right)^{2}} f \psi^{4 n \alpha}\|\nabla v\|_{\mathcal{L}}^{2}-4 n \alpha(4 n \alpha-1) f \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2}+4 n^{2} \alpha^{2} K^{\alpha} \psi^{4 n \alpha-1}
\end{aligned}
$$

Next, we compute

$$
\begin{aligned}
-4 \theta \varphi v^{-3} & \psi^{4 n \alpha}\langle\nabla v, \nabla f\rangle_{\mathcal{L}}= \\
& =-4 \theta \varphi v^{-3}\left\langle\nabla v, \nabla\left(f \psi^{4 n \alpha}\right)\right\rangle_{\mathcal{L}}+16 n \alpha \theta \varphi v^{-3} f \psi^{4 n \alpha-1}\langle\nabla v, \nabla \psi\rangle_{\mathcal{L}} \\
& \leqslant-4 \theta \varphi v^{-3}\left\langle\nabla v, \nabla\left(f \psi^{4 n \alpha}\right)\right\rangle_{\mathcal{L}}+\frac{4 \theta f \psi^{4 n \alpha}\|\nabla v\|_{\mathcal{L}}^{2}}{\left(2 \theta-v^{2}\right)^{2}}+16 n^{2} \alpha^{2} \theta\left(2 \theta-v^{2}\right)^{2} \varphi^{2} v^{-6} f \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2} \\
& =-4 \theta \varphi v^{-3}\left\langle\nabla v, \nabla\left(f \psi^{4 n \alpha}\right)\right\rangle_{\mathcal{L}}+\frac{4 \theta}{\left(2 \theta-v^{2}\right)^{2}} f \psi^{4 n \alpha}\|\nabla v\|_{\mathcal{L}}^{2}+16 n^{2} \alpha^{2} \theta v^{-2} f \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2} .
\end{aligned}
$$

Moreover, we have

$$
-2\left\langle\nabla \psi^{4 n \alpha}, \nabla f\right\rangle_{\mathcal{L}}=-2 \psi^{-4 n \alpha}\left\langle\nabla \psi^{4 n \alpha}, \nabla\left(f \psi^{4 n \alpha}\right)\right\rangle_{\mathcal{L}}+32 n^{2} \alpha^{2} f \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2} .
$$

Combining the above gives

$$
\begin{aligned}
\partial_{t}\left(f \psi^{4 n \alpha}\right) \leqslant & \mathcal{L}\left(f \psi^{4 n \alpha}\right)-\left\langle 2 \psi^{-4 n \alpha} \nabla \psi^{4 n \alpha}+4 \theta \varphi v^{-3} \nabla v, \nabla\left(f \psi^{4 n \alpha}\right)\right\rangle_{\mathcal{L}}-2 \alpha f K^{\alpha} H \varphi \psi^{4 n \alpha} \\
& +\left(32 n^{2} \alpha^{2}+16 n^{2} \alpha^{2} \theta v^{-2}-4 n \alpha(4 n \alpha-1)\right) f \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2}+4 n^{2} \alpha^{2} K^{\alpha} \psi^{4 n \alpha-1}
\end{aligned}
$$

In addition, on the support of $\psi$, we have $\nabla \psi=-\nabla \bar{u}=-\nabla\left\langle F, \vec{e}_{n+1}\right\rangle$ which leads to

$$
\begin{aligned}
\|\nabla \psi\|_{\mathcal{L}}^{2} & =\left\|\nabla\left\langle F, \vec{e}_{n+1}\right\rangle\right\|_{\mathcal{L}}^{2} \leqslant \sum_{m=1}^{n+1}\left\|\nabla\left\langle F, \vec{e}_{m}\right\rangle\right\|_{\mathcal{L}}^{2}=\sum_{m+1}^{n+1} \alpha K^{\alpha} b^{i j}\left\langle F_{i}, \vec{e}_{m}\right\rangle\left\langle F_{j}, \vec{e}_{m}\right\rangle \\
& =\sum_{i=1}^{n} \sum_{j=1}^{n} \alpha K^{\alpha} b^{i j}\left(\sum_{m=1}^{n+1}\left\langle F_{i}, \vec{e}_{m}\right\rangle\left\langle F_{j}, \vec{e}_{m}\right\rangle\right)=\sum_{i=1}^{n} \sum_{j=1}^{n} \alpha K^{\alpha} b^{i j} g_{i j} \leqslant n \alpha K^{\alpha} \lambda_{\min }^{-1} \leqslant n \alpha \Lambda K^{\alpha} .
\end{aligned}
$$

Hence, $v \geqslant 1$ implies

$$
\left(32 n^{2} \alpha^{2}+16 n^{2} \alpha^{2} \theta v^{-2}-4 n \alpha(4 n \alpha-1)\right) f \psi^{4 n \alpha-2}\|\nabla \psi\|_{\mathcal{L}}^{2} \leqslant n \alpha\left(16 n^{2} \alpha^{2}(\theta+1)+4 n \alpha\right) f \psi^{4 n \alpha-2} \Lambda K^{\alpha} .
$$

Thus, by the inequalities $H \geqslant n K^{\frac{1}{n}}$ and $\varphi \geqslant 1 /(2 \theta)$, the evolution equation of $f \psi^{4 n \alpha}$ can be reduced to the following

$$
\begin{aligned}
\partial_{t}\left(f \psi^{4 n \alpha}\right) \leqslant & \mathcal{L}\left(f \psi^{4 n \alpha}\right)-\left\langle 2 \psi^{-4 n \alpha} \nabla \psi^{4 n \alpha}+4 \theta \varphi v^{-3} \nabla v, \nabla f \psi^{4 n \alpha}\right\rangle_{\mathcal{L}}-n \alpha \theta^{-1} K^{\alpha+\frac{1}{n}} f \psi^{4 n \alpha} \\
& +4 n^{2} \alpha^{2}(4 n \alpha(\theta+1)+1) \Lambda K^{\alpha} f \psi^{4 n \alpha-2}+4 n^{2} \alpha^{2} K^{\alpha} \psi^{4 n \alpha-1} .
\end{aligned}
$$

Involving $\eta=: t(1+t)^{-1}$ and $\partial \eta_{t}=(1+t)^{-2} \leqslant 1$ yields

$$
\begin{aligned}
\partial_{t}\left(\eta^{2 n \alpha} f \psi^{4 n \alpha}\right) \leqslant & \mathcal{L}\left(\eta^{2 n \alpha} f \psi^{4 n \alpha}\right)-\left\langle 2 \psi^{-4 n \alpha} \nabla \psi^{4 n \alpha}+4 \theta \varphi v^{-3} \nabla v, \nabla \eta^{2 n \alpha} f \psi^{4 n \alpha}\right\rangle_{\mathcal{L}}-n \alpha \theta^{-1} K^{\alpha+\frac{1}{n}} \eta^{2 n \alpha} f \psi^{4 n \alpha} \\
& +4 n^{2} \alpha^{2}(4 n \alpha(\theta+1)+1) \Lambda K^{\alpha} \eta^{2 n \alpha} f \psi^{4 n \alpha-2}+4 n^{2} \alpha^{2} K^{\alpha} \eta^{2 n \alpha} f \psi^{4 n \alpha-1}+2 n \alpha \eta^{2 n \alpha-1} f \psi^{4 n \alpha} .
\end{aligned}
$$

Since $\psi$ is compactly supported, $\eta^{2 n \alpha} f \psi^{4 n \alpha}$ attains its maximum in $M^{n} \times\left[0, T_{0}\right]$ at some $\left(p_{0}, t_{0}\right)$ with $t_{0}>0$. Then, the last inequality implies that at $\left(p_{0}, t_{0}\right)$

$$
\begin{aligned}
n \alpha \theta^{-1} K^{\alpha+\frac{1}{n}} \eta^{2 n \alpha} f \psi^{4 n \alpha} \leqslant & 4 n^{2} \alpha^{2}(4 n \alpha(\theta+1)+1) \Lambda K^{\alpha} \eta^{2 n \alpha} f \psi^{4 n \alpha-2} \\
& +4 n^{2} \alpha^{2} K^{\alpha} \eta^{2 n \alpha} f \psi^{4 n \alpha-1}+2 n \alpha \eta^{2 n \alpha-1} f \psi^{4 n \alpha}
\end{aligned}
$$

Multiplying by $(n \alpha)^{-1} \theta K^{-\alpha} \eta^{-2 n \alpha+1} f^{-1} \psi^{-4 n \alpha+2}$ yields the bound

$$
\eta K^{\frac{1}{n}} \psi^{2} \leqslant 4 n \alpha \theta(4 n \alpha(\theta+1)+1) \Lambda \eta+4 n \alpha \theta \eta \psi+2 \theta K^{-\alpha} \psi^{2}
$$

and by $\theta \geqslant 1, \psi \leqslant M, 1 \leqslant M$, and $\eta \leqslant 1$

$$
\begin{aligned}
\eta K^{\frac{1}{n}} \psi^{2} & \leqslant 4 n \alpha \theta \eta((4 n \alpha(\theta+\theta)+\theta) \Lambda+\psi)+2 \theta \eta^{n \alpha} \psi^{2+2 n \alpha}\left(\eta K^{\frac{1}{n}} \psi^{2}\right)^{-n \alpha} \\
& \leqslant 4 n \alpha \theta(8 n \alpha+1)(\theta \Lambda+M)+2 \theta M^{2+2 n \alpha}\left(\eta K^{\frac{1}{n}} \psi^{2}\right)^{-n \alpha} \\
& \leqslant 2 \theta\left(16 n^{2} \alpha^{2}+2 n \alpha+M^{2 n \alpha}\left(\eta K^{\frac{1}{n}} \psi^{2}\right)^{-n \alpha}\right)\left(\theta \Lambda+M^{2}\right)
\end{aligned}
$$

Hence, in the case of $\eta K^{\frac{1}{n}} \psi^{2} \geqslant M^{2}$, the last inequality yields

$$
\eta K^{\frac{1}{n}} \psi^{2} \leqslant 2 \theta\left(16 n^{2} \alpha^{2}+2 n \alpha+1\right)\left(\theta \Lambda+M^{2}\right) \leqslant 2 \theta(4 n \alpha+1)^{2}\left(\theta \Lambda+M^{2}\right)
$$

In the other case, we can simply obtain $\eta K^{\frac{1}{n}} \psi^{2} \leqslant M^{2} \leqslant 2 \theta(4 n \alpha+1)^{2}\left(\theta \Lambda+M^{2}\right)$. Thus, at $\left(p_{0}, t_{0}\right)$,

$$
\eta K^{\frac{1}{n}} \psi^{2} \leqslant 2 \theta(4 n \alpha+1)^{2}\left(\theta \Lambda+M^{2}\right)
$$

Let $\Psi$ denote the maximum value $\eta^{2 n \alpha} f \psi^{4 n \alpha}\left(p_{0}, t_{0}\right)=\eta^{2 n \alpha} \varphi K^{2 \alpha} \psi^{4 n \alpha}\left(p_{0}, t_{0}\right)$. Then, $\varphi \leqslant 1$ gives

$$
\Psi \leqslant\left(\eta K^{\frac{1}{n}} \psi^{2}\right)^{2 \alpha n}\left(p_{0}, t_{0}\right) \leqslant(2 \theta)^{2 \alpha n}(4 n \alpha+1)^{4 \alpha n}\left(\theta \Lambda+M^{2}\right)^{2 \alpha n} .
$$

Using also that $(2 \theta)^{-1} \leqslant \varphi$, we finally conclude that for all $p \in M^{n}$ and $t \in\left[0, T_{0}\right]$ the following holds

$$
\frac{\eta^{2 n \alpha} K^{2 \alpha} \psi^{4 n \alpha}(p, t)}{2 \theta} \leqslant \eta^{2 n \alpha} \varphi K^{2 \alpha} \psi^{4 n \alpha}(p, t) \leqslant \Psi \leqslant(2 \theta)^{2 \alpha n}(4 n \alpha+1)^{4 \alpha n}(\theta \Lambda+M)^{2 \alpha n}
$$

Hence, setting $t=T_{0}$ yields

$$
\left(\eta K^{\frac{1}{n}} \psi^{2}\right)\left(p, T_{0}\right) \leqslant(2 \theta)^{1+\frac{1}{2 n \alpha}}(4 n \alpha+1)^{2}\left(\theta \Lambda+M^{2}\right)
$$

and the desired result simply follows by substituting $T_{0}$ by $t$.

## Chapter 4

## Optimal regularity of translators

A solution $\Sigma_{t}$ to the Gauss curvature flow in $\mathbb{R}^{n+1}$ translating along a direction $\vec{e}_{n+1}$ satisfies

$$
\Sigma_{t}=\Sigma+c t \vec{e}_{n+1}=:\left\{Y+c t \vec{e}_{n+1} \in \mathbb{R}^{n+1}: Y \in \Sigma\right\}
$$

where the speed $c$ is a constant, and $\Sigma$ is a complete convex hypersurface embedded in $\mathbb{R}^{n+1}$. We observe that there exist a convex open set $\Omega \subset \mathbb{R}^{n}$ and a convex function $u: \Omega \rightarrow \mathbb{R}$ satisfying

$$
\Sigma \text { is the boundary of }\{(x, t): x \in \Omega, t \geqslant u(x)\} .
$$

By the result in [27], the set $\Omega$ must be bounded. If $\mathcal{A}(\Omega)$ denotes the area of $\Omega$, it follows that $u$ is a smooth function satisfying

$$
\left\{\begin{array}{l}
\frac{\operatorname{det} D^{2} u}{\left(1+|D u|^{2}\right)^{\frac{n+1}{2}}}=\frac{\left|S^{n-1}\right|}{\mathcal{A}(\Omega)} \text { in } \Omega,  \tag{4.1}\\
\lim _{x \rightarrow \partial \Omega}|D u|(x)=+\infty \quad \text { on } \partial \Omega .
\end{array}\right.
$$

where $S^{n}$ is the unit $n$-sphere.
Conversely, given an open bounded convex set $\Omega \subset \mathbb{R}^{n}$, there exists a solution $u: \Omega \rightarrow \mathbb{R}$ of (4.1), and any two solutions differ by a constant. (See [27] and Theorem 4.8 in [26]). Hence, given a translator $\Sigma$ in $\mathbb{R}^{n+1}$ of the Gauss curvature flow, there exists an open bounded convex set $\Omega \in \mathbb{R}^{n}$ such that $\Sigma$ converges to the cylinder $\partial \Omega \times \mathbb{R}$, and the immersion $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ of $F\left(M^{n}\right)=\Sigma$ satisfies

$$
\begin{equation*}
K(p)=\frac{\left|S^{n-1}\right|}{\mathcal{A}(\Omega)}\left\langle\vec{n}(p), \vec{e}_{n+1}\right\rangle . \tag{4.2}
\end{equation*}
$$

We recall the result of John Urbas in [27].
Theorem 4.1 (Urbas). Given an open bounded convex domain $\Omega \subset \mathbb{R}^{2}$, there exists a convex solution $u: \Omega \rightarrow \mathbb{R}$ satisfying (4.1), and it is unique up to addition by a constant. In particular,
if for each $x_{0} \in \partial \Omega$, there exists a ball $B \subset \mathbb{R}^{2}$ satisfying $\Omega \subset B$ and $x_{0} \in \partial B$, then the solution $u$ is a smooth function satisfying

$$
\lim _{x \rightarrow \partial \Omega} u(x)=+\infty .
$$

This result guarantees that there exists a unique $C^{1}$ translator $\Sigma=\partial\{(x, t): x \in \Omega, t \geqslant u(x)\}$ for any open bounded convex domain $\Omega$. Also, if $\Omega$ is a uniformly convex domain, then $\Sigma$ is strictly convex, and thus $C^{\infty}$ smooth by standard estimates. However, if $\Omega$ is weakly convex, then $\Sigma$ may not be strictly convex on the boundary of $\Omega$. Richard Hamilton conjectured that if $\Omega$ is a square, then $\Sigma$ has flat sides on the boundary of $\Omega$. This is shown in the next picture.


Figure 1. Translator $\Sigma$ on a square

The Hamilton's conjecture and the optimal regularity of the translator was shown in [11].
Theorem 4.2 (Choi-Daskalopoulos-Lee). Let $\Omega:=(-1,1) \times(-1,1) \subset \mathbb{R}^{2}$ be the open square and $V$ be the set of the vertices $( \pm 1, \pm 1)$ of $\partial \Omega$. Assume $u: \Omega \rightarrow \mathbb{R}$ is a convex smooth solution of (4.1), and let $\Sigma$ denote the boundary of $\{(x, t): x \in \Omega, t \geqslant u(x)\}$. Then, $\Sigma$ is a complete convex hypersurface of class $C_{\text {loc }}^{1,1}$, and there exists a smooth function $\bar{u}:(\partial \Omega \backslash V) \rightarrow \mathbb{R}$ satisfying

$$
\lim _{x \rightarrow y} u(x)=\bar{u}(y), \quad \quad \lim _{x \rightarrow V} u(x)=\lim _{y \rightarrow V} \bar{u}(y)=+\infty
$$

Remark 4.3 (Optimal regularity). Let us show that if there exists a flat side on a translator $\Sigma$, then $\Sigma$ has at most $C^{1,1}$ regularity. Assume that the graph $(x, h(x, z), z)$ of a function $h(x, z)$ is a part of a translator. Then, we have

$$
h_{v v} h_{\tau \tau} \geqslant \frac{\operatorname{det} D^{2} h}{\left(1+|D h|^{2}\right)^{\frac{3}{2}}} \geqslant-\frac{2 \pi}{\mathcal{A}}|D h|\left\langle v, e_{2}\right\rangle
$$

where $v(x, z)$ is the outward normal and $\tau(x, z)$ is a tangential direction of the level set of $h(x, z)$ at a point $(x, z)$. We denote by $L_{r}$ the level set $\{(x, z): h(x, z)=r\}$, and denote by $\kappa(x, z)$ the curvature of $L_{h(x, z)}$ at $(x, z)$. Since we have $h_{\tau \tau}=|D h| \kappa$, the inequality above gives $h_{v v} \kappa \geqslant-2 \pi \mathcal{A}^{-1}\left\langle v, e_{2}\right\rangle$. We will establish the local lower bound for $-\left\langle v, e_{2}\right\rangle$ in section 3 , which guarantees that

$$
h_{v v} K \geqslant c .
$$

We choose a neighborhood $U$ of a point $\left(x_{0}, z_{0}\right)$ on the free boundary $\Gamma$, namely $\left(x_{0}, z_{0}\right) \in \Gamma=$ : $\partial L_{-1}$. Since the level sets $L_{r}$ monotonically converge to $\Gamma$, there exists a constant $c$ such that $\int_{L_{r} \cap U} d s \geqslant c$ for $r$ close enough to -1 , where $s$ is the arc length parameter. Hence, the following holds

$$
2 \pi \geqslant \int_{L_{r}} \kappa d s \geqslant \int_{L_{r} \cap U} \kappa d s \geqslant \frac{c}{\max _{L_{r} \cap U} h_{v v}} .
$$

Thus, $\max _{L_{r} \cap U} h_{\nu v} \geqslant c$ holds for some uniform constant $c$. However, we have $D^{2} h=0$ on $L_{-1}$. Therefore, $D^{2} h$ is not a continuous function.

## 1 Optimal $C^{1,1}$ regularity

In this section, we will establish a local curvature estimate for smooth strictly convex complete solutions of equation (4.2). In the last section we will use this estimate to obtain the optimal $C^{1,1}$ regularity for a weakly convex solution of (4.2) in the degenerate case. We recall that a solution of (4.2) has an immersion $F: M^{2} \rightarrow \mathbb{R}^{3}$ of $F\left(M^{2}\right)=\Sigma$. Given a ball $B_{R}(Y)$ we define the associated cut-off function $\eta$ by $\eta(p)=\left(|F(p)-Y|^{2}-R^{2}\right)_{+}$. We have the following result.

Theorem 4.4 (Curvature bound). Let $\Sigma$ be a smooth strictly convex complete solution of (4.2). Let $\Sigma^{c}$ be the cut off from $\Sigma$ by a ball $B_{R}(Y) \subset \mathbb{R}^{n+1}$. Then, for any $p \in M^{2}$ with $F(p) \in \Sigma^{c}$, the maximum principal curvature $\Lambda(p):=\max \left\{\lambda_{1}(p), \lambda_{2}(p)\right\}$ satisfies

$$
\eta \Lambda(p) \leqslant \frac{9 \pi}{\mathcal{A}(\Omega)} \sup _{F(q) \in \Sigma^{c}}|F(q)-Y|^{3} .
$$

Proof. We may assume, without loss of generality, that $Y=0$. Recall the definition of the cutting ball. The continuous function $\eta \Lambda$ attains its maximum on the compact set $\Sigma^{c}$ at some point $F\left(p_{0}\right) \in \Sigma^{c}$,

$$
\eta \Lambda\left(p_{0}\right)=\max _{F(p) \in \Sigma^{c}} \eta \Lambda(p) .
$$

Then, because we have $\eta=0$ on $\partial \Sigma_{c}, F\left(p_{0}\right)$ is an interior point of $\Sigma_{c}$. Thus, $\eta \Lambda$ attains a local maximum at $p_{0}$. Moreover, we can choose an open chart $(U, \varphi)$ with $p_{0} \in \varphi(U)$ and $F(\varphi(U)) \subset \Sigma_{c}$ such that the covariant derivatives $\left\{\nabla_{1} F\left(p_{0}\right), \nabla_{2} F\left(p_{0}\right)\right\}$ form an orthonormal basis of $T \Sigma_{F\left(p_{0}\right)}$ satisfying

$$
g_{i j}\left(p_{0}\right)=\delta_{i j}, \quad h_{i j}\left(p_{0}\right)=\delta_{i j} \lambda_{i}\left(p_{0}\right), \quad \lambda_{1}\left(p_{0}\right)=\Lambda\left(p_{0}\right)
$$

Next, we define the function $w: U \rightarrow \mathbb{R}$ by

$$
w=\eta \frac{h_{11}}{g_{11}} .
$$

Then, the Euler formula, Proposition 2.3 guarantees $w \leqslant \eta \Lambda$. Therefore, for all $p \in U$, the following holds

$$
w(p) \leqslant \eta \Lambda(p) \leqslant \eta \Lambda\left(p_{0}\right)=w\left(p_{0}\right) .
$$

Thus, $w$ also attains its maximum at $p_{0}$.
Now, we consider the derivative of $w$. Then, $\nabla g_{11}=0$ gives

$$
\begin{equation*}
\frac{\nabla_{i} w}{w}=\frac{\nabla_{i} h_{11}}{h_{11}}+\frac{\nabla_{i} \eta}{\eta} . \tag{4.3}
\end{equation*}
$$

Differentiating the equation above yields

$$
\frac{\nabla_{i} \nabla_{j} w}{w}-\frac{\nabla_{i} w \nabla_{j} w}{w^{2}}=\frac{\nabla_{i} \nabla_{j} h_{11}}{h_{11}}-\frac{\nabla_{i} h_{11} \nabla_{j} h_{11}}{\left(h_{11}\right)^{2}}+\frac{\nabla_{i} \nabla_{j} \eta}{\eta}-\frac{\nabla_{i} \eta \nabla_{j} \eta}{\eta^{2}}
$$

and multiplying by $K b^{i j}$, we obtain

$$
\begin{equation*}
\frac{\mathcal{L} w}{w}-\frac{\|\nabla w\|_{\mathcal{L}}^{2}}{w^{2}}=\frac{\mathcal{L} h_{11}}{h_{11}}-\frac{\left\|\nabla h_{11}\right\|_{\mathcal{L}}^{2}}{\left(h_{11}\right)^{2}}+\frac{\mathcal{L} \eta}{\eta}-\frac{\|\nabla \eta\|_{\mathcal{L}}^{2}}{\eta^{2}} . \tag{4.4}
\end{equation*}
$$

Observing next that $\mathcal{L} F:=K b^{i j} \nabla_{j} \nabla_{j} F=K b^{i j} h_{i j} \vec{n}=2 K \vec{n}$, we compute $\mathcal{L} \eta$ on the support of $\eta$ as follows:

$$
\mathcal{L} \eta=\mathcal{L}|F|^{2}=2\langle F, \mathcal{L} F\rangle+2\langle\nabla F, \nabla F\rangle_{\mathcal{L}}=4 K\langle F, \vec{n}\rangle+2 K b^{i j} g_{i j}=4 K\langle F, \vec{n}\rangle+2 H .
$$

Thus, $4 K=8 \pi \mathcal{A}^{-1}\left\langle\vec{e}_{3}, \vec{n}\right\rangle \leqslant 8 \pi \mathcal{A}^{-1}$ and $2 H \geqslant 2 \Lambda$ imply

$$
\begin{equation*}
\mathcal{L} \eta \geqslant-8 \pi \mathcal{A}^{-1}|F|+2 \Lambda . \tag{4.5}
\end{equation*}
$$

Since $w$ attains its maximum at $p_{0}$, we have $\nabla w\left(p_{0}\right)=0$, and thus (4.3) gives

$$
\frac{\left\|\nabla h_{11}\right\|_{\mathcal{L}}^{2}}{\left(h_{11}\right)^{2}}\left(p_{0}\right)=\frac{\|\nabla \eta\|_{\mathcal{L}}^{2}}{\eta^{2}}\left(p_{0}\right)
$$

Hence, combining $\mathcal{L} w\left(p_{0}\right) \leqslant 0$, (4.4), (4.5) and the equation above yields the following at $p_{0}$

$$
\begin{equation*}
0 \geqslant \frac{\mathcal{L} h_{11}}{h_{11}}+\frac{2 \Lambda}{\eta}-\frac{8 \pi|F|}{\mathcal{A} \eta}-\frac{2\left\|\nabla h_{11}\right\|_{\mathcal{L}}^{2}}{\left(h_{11}\right)^{2}} . \tag{4.6}
\end{equation*}
$$

To compute $\mathcal{L} h_{11}$, we begin by differentiating $K$,

$$
\begin{equation*}
\nabla_{1} K=K b^{i j} \nabla_{1} h_{i j} . \tag{4.7}
\end{equation*}
$$

By differentiating the equation above again, we obtain

$$
\begin{equation*}
\nabla_{1} \nabla_{1} K=K b^{i j} \nabla_{1} \nabla_{1} h_{i j}+K b^{i j} b^{k l} \nabla_{1} h_{i j} \nabla_{1} h_{k l}-K b^{i k} b^{j l} \nabla_{1} h_{i j} \nabla_{1} h_{k l} . \tag{4.8}
\end{equation*}
$$

We can derive $\mathcal{L} h_{11}$ from the first term $K b^{i j} \nabla_{1} \nabla_{1} h_{i j}$ as follows

$$
\begin{align*}
K b^{i j} \nabla_{1} \nabla_{1} h_{i j} & =K b^{i j} \nabla_{1} \nabla_{i} h_{j 1}=K b^{i j}\left(\nabla_{i} \nabla_{1} h_{j 1}+R_{1 i j k} h_{1}^{k}+R_{1 i 1 k} h_{j}^{k}\right)  \tag{4.9}\\
& =K b^{i j} \nabla_{i} \nabla_{j} h_{11}+K b^{i j}\left(h_{1 j} h_{i k}-h_{1 k} h_{i j}\right) h_{1}^{k}+K b^{i j}\left(h_{11} h_{i k}-h_{1 k} h_{i 1}\right) h_{j}^{k} \\
& =\mathcal{L} h_{11}-2 K h_{1 k} h_{1}^{k}+K H h_{11} .
\end{align*}
$$

On the other hand, differentiating (4.2) yields

$$
\begin{equation*}
\nabla_{1} K=\frac{2 \pi}{\mathcal{A}}\left\langle\nabla_{1} \vec{n}, \vec{e}_{3}\right\rangle=-\frac{2 \pi}{\mathcal{A}} h_{1 k}\left\langle F^{k}, \vec{e}_{3}\right\rangle \tag{4.10}
\end{equation*}
$$

To get the right hand side of (4.8), we differentiate the equation above,

$$
\begin{equation*}
\nabla_{1} \nabla_{1} K=-\frac{2 \pi}{\mathcal{A}} \nabla_{1} h_{1 k}\left\langle F^{k}, \vec{e}_{3}\right\rangle-\frac{2 \pi}{\mathcal{A}} h_{1 k} h_{1}^{k}\left\langle\vec{n}, \vec{e}_{3}\right\rangle=-\frac{2 \pi}{\mathcal{A}} \nabla_{k} h_{11}\left\langle F^{k}, \vec{e}_{3}\right\rangle-K h_{1 k} h_{1}^{k} \tag{4.11}
\end{equation*}
$$

Combining (4.8), (4.9), and (4.11), we obtain the following at $p_{0}$

$$
\mathcal{L} h_{11}=2\left|\nabla_{2} h_{11}\right|^{2}-2 \nabla_{1} h_{11} \nabla_{1} h_{22}-\frac{2 \pi}{\mathcal{A}} \nabla_{k} h_{11}\left\langle F^{k}, \vec{e}_{3}\right\rangle-K^{2} .
$$

Hence, at $p_{0}$, applying the equation above to (4.6) and the definition of the norm $\|\cdot\|_{\mathcal{L}}^{2}$ yield

$$
\begin{aligned}
0 \geqslant & \frac{1}{h_{11}}\left(2\left|\nabla_{2} h_{11}\right|^{2}-2 \nabla_{1} h_{11} \nabla_{1} h_{22}-\frac{2 \pi}{\mathcal{A}} \nabla_{k} h_{11}\left\langle F^{k}, \vec{e}_{3}\right\rangle-K^{2}\right) \\
& -\frac{2 h_{22}\left|\nabla_{1} h_{11}\right|^{2}+2 h_{11}\left|\nabla_{2} h_{11}\right|^{2}}{\left(h_{11}\right)^{2}}+\frac{2 \Lambda}{\eta}-\frac{8 \pi|F|}{\mathcal{A} \eta} \\
= & -\frac{2 \nabla_{1} h_{11}\left(h_{22} \nabla_{1} h_{11}+h_{11} \nabla_{1} h_{22}\right)}{\left(h_{11}\right)^{2}}+\frac{1}{h_{11}}\left(-\frac{2 \pi}{\mathcal{A}} \nabla_{k} h_{11}\left\langle F^{k}, \vec{e}_{3}\right\rangle-K^{2}\right)+\frac{2 \Lambda}{\eta}-\frac{8 \pi|F|}{\mathcal{A} \eta} .
\end{aligned}
$$

However, (4.7) and (4.10) imply the following at $p_{0}$

$$
h_{22} \nabla_{1} h_{11}+h_{11} \nabla_{1} h_{22}=\nabla_{1} K=-\frac{2 \pi}{\mathcal{A}} h_{11}\left\langle F^{1}, \vec{e}_{3}\right\rangle .
$$

Therefore, the last inequality can be reduced to

$$
0 \geqslant-\frac{2 \pi}{\mathcal{A}} \frac{\nabla_{2} h_{11}}{h_{11}}\left\langle F^{2}, \vec{e}_{3}\right\rangle+\frac{2 \pi}{\mathcal{A}} \frac{\nabla_{1} h_{11}}{h_{11}}\left\langle F^{1}, \vec{e}_{3}\right\rangle-\frac{K^{2}}{h_{11}}+\frac{2 \Lambda}{\eta}-\frac{8 \pi|F|}{\mathcal{A} \eta} .
$$

Observing $\left(h_{11}\right)^{-1} \nabla_{i} h_{11}\left(p_{0}\right)=-\eta^{-1} \nabla_{i} \eta\left(p_{0}\right)$ by (4.3) and $\nabla w\left(p_{0}\right)=0$, we have

$$
0 \geqslant \frac{2 \pi}{\mathcal{A}} \frac{\nabla_{2} \eta}{\eta}\left\langle F^{2}, \vec{e}_{3}\right\rangle-\frac{2 \pi}{\mathcal{A}} \frac{\nabla_{1} \eta}{\eta}\left\langle F^{1}, \vec{e}_{3}\right\rangle-\frac{K^{2}}{h_{11}}+\frac{2 \Lambda}{\eta}-\frac{8 \pi|F|}{\mathcal{A} \eta} .
$$

Applying $h_{11}\left(p_{0}\right)=\Lambda\left(p_{0}\right)$ and (4.2) to the inequality above, we obtain

$$
0 \geqslant \frac{4 \pi}{\mathcal{A} \eta}\left\langle F_{2}, F\right\rangle\left\langle F^{2}, \vec{e}_{3}\right\rangle-\frac{4 \pi}{\mathcal{A} \eta}\left\langle F_{1}, F\right\rangle\left\langle F^{1}, \vec{e}_{3}\right\rangle-\frac{4 \pi^{2}\left|\left\langle\vec{n}, \vec{e}_{3}\right\rangle\right|^{2}}{\mathcal{A}^{2} \Lambda}+\frac{2 \Lambda}{\eta}-\frac{8 \pi|F|}{\mathcal{A} \eta} .
$$

We next multiply by $\eta \Lambda$ the last inequality and apply $\left|\left\langle F_{i}, F\right\rangle\left\langle F^{i}, \vec{e}_{3}\right\rangle\right| \leqslant|F|$ and $\left\langle\vec{n}, \vec{e}_{3}\right\rangle \leqslant 1$. Then, by also using the definition $\eta=:\left(|F|^{2}-R^{2}\right)$, we obtain

$$
0 \geqslant-\frac{16 \pi}{\mathcal{A}}|F| \Lambda-\frac{4 \pi^{2} \eta}{\mathcal{A}^{2}}+2 \Lambda^{2} \geqslant 2 \Lambda^{2}-\frac{16 \pi}{\mathcal{A}}|F| \Lambda-\frac{4 \pi^{2}}{\mathcal{A}^{2}}|F|^{2} .
$$

Solving the quadratic inequality of $\Lambda$, we obtain an upper bound of $\Lambda$ at $p_{0}$,

$$
\Lambda \leqslant \frac{(4+3 \sqrt{2}) \pi|F|}{\mathcal{A}} \leqslant \frac{9 \pi}{\mathcal{A}}|F| .
$$

Therefore, multiplying by $\eta \leqslant|F|^{2}$ yields the desired result,

$$
\eta \Lambda(p) \leqslant \eta \Lambda\left(p_{0}\right) \leqslant \frac{9 \pi}{\mathcal{A}}|F|^{3}\left(p_{0}\right) \leqslant \frac{9 \pi}{\mathcal{A}} \sup _{F(q) \in \Sigma_{c}}|F|^{3}(q) .
$$

## 2 Partial derivative bound

Definition 4.5 (Axial symmetry). We say that a surface $\Sigma \subset \mathbb{R}^{3}$ has axial symmetry, if $(x, y, z) \in \Sigma$ guarantees $(-x, y, z),(x,-y, z) \in \Sigma$. Similarly, a set $\Omega \subset \mathbb{R}^{2}$ has axial symmetry, if $(x, y) \in \Omega$ guarantees $(-x, y),(x,-y) \in \Omega$.

Notation 4.6 (To be used in sections 4 and 5). Also, we summarize some further notation.
(i) Given a set $A \subset \mathbb{R}^{3}$ and a constant $s$, we denote the $x=s$ level set by $L_{s}^{x}(A)=$ $\{(s, y, z) \in A\}$. Similarly, we denote $y=s$ and $z=s$ level set by $L_{s}^{y}(A)$ and $L_{s}^{z}(A)$, respectively.
(ii) Given a constant $s$ and a function $f: \Omega \rightarrow \mathbb{R}$ with $\Omega \subset \mathbb{R}^{2}$, we denote by $L_{s}(f)$ the $s$-level set $\{(x, y) \in \Omega: f(x, y)=s\}$.
(iii) We let $e_{1}$ and $e_{2}$ the unit vectors $(1,0)$ and $(0,1)$, respectively.
(iv) For a complete and convex curve $\Gamma \subset \mathbb{R}^{2}$, its convex hull $\operatorname{Conv}(\Gamma)$ is given by

$$
\operatorname{Conv}(\Gamma)=\{(t x+(1-t) y: x, y \in \Gamma, t \in[0,1]\} .
$$

If $A$ is a subset of $\operatorname{Conv}(\Gamma)$, then we say $A$ is enclosed by $\Gamma$ and use the notation

$$
A<\Gamma
$$

(v) Given a set $A \subset \mathbb{R}^{2}, \operatorname{cl}(A)$ and $\operatorname{Int}(A)$ mean the closure and the interior of $A$, respectively.

Assume that $\Omega$ is an open bounded strictly convex and smooth domain of $\mathbb{R}^{2}$ and assume in addition that $\Omega$ is axially symmetric. Let $u$ be the unique solution of (4.1) on $\Omega$ which defines the surface $\Sigma$. The following simple property readily follows from the uniqueness of solutions.

Proposition 4.7 (Symmetry of solutions). Let $\Omega$ be an open bounded strictly convex and smooth subset of $\mathbb{R}^{2}$ which is axially symmetric. Then, a solution $u$ of (4.1) also is axially symmetric.

The symmetry of $u$ implies that the half surface $\{(x, y, z) \in \Sigma: y \leqslant 0\}$ is the graph of a function $h: \Omega_{y} \rightarrow \mathbb{R}$, that is

$$
\{(x, y, z) \in \Sigma: y \leqslant 0\}=\left\{(x, h(x, z), z):(x, z) \in \Omega_{y}\right\}, \quad \text { where } \quad \Omega_{y}=\{(x, z):(x, y, z) \in \Sigma\} .
$$

Moreover, the function $h$ satisfies the following equation

$$
\begin{equation*}
\frac{\operatorname{det} D^{2} h}{\left(1+|D h|^{2}\right)^{\frac{3}{2}}}=K\left(1+|D h|^{2}\right)^{\frac{1}{2}}=\frac{2 \pi}{\mathcal{A}(\Omega)}\left\langle\vec{e}_{3}, \vec{n}\right\rangle\left(1+|D h|^{2}\right)^{\frac{1}{2}}=-\frac{2 \pi}{\mathcal{A}(\Omega)} h_{z} . \tag{4.12}
\end{equation*}
$$

The right hand side of the equation above can written as $(2 \pi / \mathcal{A}) h_{v}\left\langle-e_{2}, v\right\rangle$, where $v$ is the outward normal direction of the level set of $h$. Thus, the degenerate Monge-Ampere equation (4.12) has two degenerating factors $h_{v}=|D h|$ and $\left\langle-e_{2}, v\right\rangle$. In this section, we study the lower bound for $\left\langle-e_{2}, v\right\rangle$ which corresponds to the upper bound of $\left|\partial_{x} u\right|$, the partial derivative bound. Notice that $\left|\partial_{x} u\right|$ is bounded even on the flat sides.

To obtain the lower bound for $\left\langle-e_{2}, v\right\rangle$, we will construct an one parameter family of very wide but short supersolutions $\varphi_{\alpha}$ of equation (4.12) with $\mathcal{A}(\Omega)<8$. We will then cut the graph of $\varphi_{\alpha}$ so that each of them each contained in a narrow cylinder. By sliding $\varphi_{\alpha}$ along the $z$-axis we will estimate the partial derivative $\partial_{x} u$ at a touching point which will lead to the desired upper bound for $\partial_{x} u$ as stated in Theorem 4.10.

Definition 4.8 (Barrier construction). Given a constant $\alpha \in(0,1 / 6)$, denote by $\Delta_{\alpha}$ the convex set
$\Delta_{\alpha}=\left\{(x, z) \in \mathbb{R}^{2}: x \in\left(-\frac{2}{\alpha}+1-3 \alpha, 1-3 \alpha\right), z \geqslant-\frac{2}{\alpha \pi} \log \cos \left(\frac{\alpha \pi}{2}\left(x-1+3 \alpha+\frac{1}{\alpha}\right)\right)\right\}$.


Figure 2. Supersolution $\varphi_{\alpha}$
We denote by $d_{\Delta_{\alpha}}(x, z)$ the distance function $d\left((x, z), \Delta_{\alpha}\right)$. In particular, if $(x, z) \in \Delta_{\alpha}$, then $d_{\Delta_{\alpha}}(x, z)=0$. By using $d_{\Delta_{\alpha}}(x, z)$, we define the $2 \alpha$-extension $\left(\Delta_{\alpha}\right)^{2 \alpha}$ of $\Delta_{\alpha}$ by

$$
\left(\Delta_{\alpha}\right)^{2 \alpha}=\left\{(x, z) \in \mathbb{R}^{2}: d_{\Delta_{\alpha}}(x, z) \leqslant 2 \alpha\right\} .
$$

Finally, we define the function $\varphi_{\alpha}: \operatorname{cl}\left(\left(\Delta_{\alpha}\right)^{2 \alpha} \backslash \Delta_{\alpha}\right) \rightarrow \mathbb{R}$ by

$$
\varphi_{\alpha}(x, z)=-1+2 \alpha-\sqrt{4 \alpha^{2}-d^{2}\left(\Delta_{\alpha}\right)(x, z)} .
$$

This is all shown in Figure 3.
Lemma 4.9 (Supersolution). Given a constant $\alpha \in(0,1 / 6)$, the function $\varphi_{\alpha}$ in Definition 4.8 is a convex function satisfying

$$
\frac{\operatorname{det} D^{2} \varphi}{\left(1+|D \varphi|^{2}\right)^{\frac{3}{2}}} \leqslant-\frac{\pi}{4} \varphi_{z} .
$$

Proof. For convenience, we let $\varphi$ and $d$ denote $\varphi_{\alpha}$ and $d_{\Delta_{\alpha}}$ respectively. For each point $p \in \mathbb{R}^{2}$ with $d(p)>0$, we denote by $\tau(p)$ and $v(p)$ the tangential and the normal direction of a level set $L_{d(p)}(d)$ of the distance function $d$ satisfying $\left\langle\tau, e_{1}\right\rangle \geqslant 0$ and $\left\langle v, e_{2}\right\rangle \leqslant 0$, respectively. Then, we have

$$
\begin{equation*}
D d=v, \quad d_{v}=|D d|=1, \quad d_{\tau}=0 . \tag{4.13}
\end{equation*}
$$

We observe $v(p)=v(p+\epsilon v(p))$ for all $\epsilon \in \mathbb{R}$ with $d(p+\epsilon \nu(p))>0$, which implies

$$
\begin{equation*}
d_{v v}=d_{v \tau}=0 . \tag{4.14}
\end{equation*}
$$

To derive $d_{\tau \tau}(p)$, given a point $p_{0}$, we consider the immersion $\gamma: \mathbb{R} \rightarrow \mathbb{R}^{2}$ satisfying $d(\gamma(s))=$ $d\left(p_{0}\right)$ with $\gamma(0)=p_{0}$, where $s$ is the arc length parameter of the level set, $\gamma(\mathbb{R})=L_{d\left(p_{0}\right)}(d)$. By differentiating $d(\gamma(s))=d\left(p_{0}\right)$ twice with respect to $s$, we obtain

$$
\left\langle\gamma_{s},\left(D^{2} d\right) \gamma_{s}\right\rangle+\left\langle D d, \gamma_{s s}\right\rangle=0
$$

We can observe that $\gamma_{s}(0)=\tau\left(p_{0}\right)$ and $\gamma_{s s}(0)=-\kappa\left(p_{0}\right) v\left(p_{0}\right)$, where $\kappa(p)>0$ is the curvature of $L_{d(p)}(d)$ at $p$. Hence, $D d(\gamma(0))=v\left(p_{0}\right)$ implies $\left\langle\tau,\left(D^{2} d\right) \tau\right\rangle+\langle-\kappa v, v\rangle=0$ at $p_{0}$. Thus,

$$
\begin{equation*}
d_{\tau \tau}(p)=\kappa(p) . \tag{4.15}
\end{equation*}
$$

Hence we can directly derive from (4.13). (4.14) and (4.15) the following holding at each $p \in$ $\operatorname{cl}\left(\left(\Delta_{\alpha}\right)^{2 \alpha} \backslash \Delta_{\alpha}\right)$

$$
\varphi_{v}=d\left(4 \alpha^{2}-d^{2}\right)^{-\frac{1}{2}}, \quad \varphi_{\tau}=0, \quad \varphi_{v v}=4 \alpha^{2}\left(4 \alpha^{2}-d^{2}\right)^{-\frac{3}{2}}, \quad \varphi_{v \tau}=0, \quad \varphi_{\tau \tau}=\kappa \varphi_{v} .
$$

Therefore, $\varphi$ is a convex function.
Next, combining the equalities above yields

$$
\begin{equation*}
\frac{\operatorname{det} D^{2} \varphi}{\left(1+|D \varphi|^{2}\right)^{\frac{3}{2}}}=\frac{\varphi_{v v} \varphi_{\tau \tau}}{\left(1+\varphi_{v}^{2}\right)^{\frac{3}{2}}}=\frac{1}{2 \alpha} \varphi_{\tau \tau}=\frac{1}{2 \alpha} \kappa \varphi_{v} . \tag{4.16}
\end{equation*}
$$

Now, we consider the point $p_{0}=p-d(p) v(p) \in \partial \Delta_{\alpha}$. Then, the convexity of $\partial \Delta_{\alpha}$ leads to

$$
\begin{equation*}
\kappa(p) \leqslant \kappa\left(p_{0}\right) . \tag{4.17}
\end{equation*}
$$

We recall that the Grim Reaper curve $\partial \Delta_{\alpha}$ is the graph of the convex function $f_{\alpha}(x)$ defined by

$$
\begin{equation*}
f_{\alpha}(x)=-\frac{2}{\pi \alpha} \log \cos \left(\frac{\alpha \pi}{2}\left(x-1+3 \alpha+\frac{1}{\alpha}\right)\right) \tag{4.18}
\end{equation*}
$$

Hence, at $x_{0}$ with $p_{0}=\left(x_{0}, f_{\alpha}\left(x_{0}\right)\right)$, the following holds

$$
\kappa\left(p_{0}\right)=\frac{f_{\alpha}^{\prime \prime}\left(x_{0}\right)}{\left(1+\left|f_{\alpha}^{\prime}\left(x_{0}\right)\right|^{2}\right)^{\frac{3}{2}}}=\frac{\pi \alpha / 2}{\left(1+\left|f_{\alpha}^{\prime}\left(x_{0}\right)\right|^{2}\right)^{\frac{1}{2}}}=-\frac{\pi \alpha}{2}\left\langle v\left(p_{0}\right), e_{2}\right\rangle .
$$

Thus, $v\left(p_{0}\right)=v(p-d(p) v(p))=v(p)$ implies

$$
\begin{equation*}
\kappa(p) \leqslant \kappa\left(p_{0}\right)=-\frac{\pi \alpha}{2}\left\langle v\left(p_{0}\right), e_{2}\right\rangle=-\frac{\pi \alpha}{2}\left\langle v(p), e_{2}\right\rangle . \tag{4.19}
\end{equation*}
$$

Therefore, given a point $p \in \operatorname{cl}\left(\left(\Delta_{\alpha}\right)^{2 \alpha} \backslash \Delta_{\alpha}\right)$, (4.16), (4.17), and (4.19) give the desired result

$$
\frac{\operatorname{det} D^{2} \varphi}{\left(1+|D \varphi|^{2}\right)^{\frac{3}{2}}}=\frac{1}{2 \alpha} \varphi_{v} K \leqslant-\frac{\pi}{4} \varphi_{v}\left\langle v, e_{2}\right\rangle=-\frac{\pi}{4}\left(\varphi_{v}\left\langle v, e_{2}\right\rangle+\varphi_{\tau}\left\langle\tau, e_{2}\right\rangle\right)=-\frac{\pi}{4} \varphi_{z} .
$$

Theorem 4.10 (Partial derivative bound). Let $\Omega$ be an open strictly convex smooth subset of $\mathbb{R}^{2}$ which is in additionally axially symmetric and satisfies

$$
[-1,1] \times[-1,1] \subset \Omega \subset\left[-\frac{4}{3}, \frac{4}{3}\right] \times\left[-\frac{4}{3}, \frac{4}{3}\right]
$$

Then, a solution $u: \Omega \rightarrow \mathbb{R}$ of (4.1) satisfies

$$
\partial_{x} u(1-5 \alpha,-1+\epsilon) \leqslant 2+3 \cot \frac{\pi \alpha^{2}}{2}
$$

for all $\epsilon \in(0,1 / 2)$ and $\alpha \in(0,1 / 6)$.

Proof. Since the domain $\Omega$ has strictly convex and smooth boundary, Theorem 4.1 implies that the graph $\Sigma=\{(x, y, u(x, y)):(x, y) \in \Omega\}$ of the function $u$ is a strictly convex complete smooth solution of (4.2). In addition by Proposition 4.7 the function $u$ is axially symmetric and we can define a convex set $\Omega_{y}$ and a convex function $h: \Omega_{y} \rightarrow \mathbb{R}$ by

$$
\Omega_{y}=\{(x, z):(x, y, z) \in \Sigma\}, \quad\{(x, y, z) \in \Sigma: y \leqslant 0\}=\left\{(x, h(x, z), z):(x, z) \in \Omega_{y}\right\} .
$$

Then, $(x, y, u(x, y))=(x, h(x, z), z)$ implies that the function $h$ satisfies 4.12). Thus, by the given condition $\mathcal{A}(\Omega) \leqslant(8 / 3)^{2}<8$ and Lemma 4.9, the function $\varphi_{\alpha}$ in Definition 4.8 is a supersolution for (4.12).

To construct a barrier $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$, we will cut the graph of $\epsilon+\varphi_{\alpha}$ by $\{1-4 \alpha\} \times \mathbb{R}^{2}$ (the blue section in Figure $3(\mathrm{a})$, and slide it along $z$-direction until it touches $\Sigma$ at a point $P_{0}$. We will show that the contact point $P_{0}$ is contained in $\{1-4 \alpha\} \times \mathbb{R}^{2}$, namely $P_{0}$ is a point on the front part of the boundary $\partial \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ of the barrier. (See the blue curve $\Gamma_{F}$ in Figure $3(\mathrm{a})$. Then, we will estimate the partial derivative $\partial_{x} u$ at $P_{0}$ by comparing with the barrier $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ at $P_{0}$. After obtaining the bound on $\partial_{x} u$ at $P_{0}$, we will use the convexity of the solution $\Sigma$ the barrier $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ to show the desired bound of $\partial_{x} u$ at $(1-5 \alpha,-1+\epsilon)$.


## Figure 3. Sliding barrier

Step 1 : Sliding barrier construction. We denote by $\Phi_{\alpha}$ the graph of $\varphi_{\alpha}$ in $[1-4 \alpha,+\infty) \times \mathbb{R}^{2}$,

$$
\Phi_{\alpha}=\left\{\left(x, \varphi_{\alpha}(x, z), z\right):(x, z) \in \operatorname{cl}\left(\left(\Delta_{\alpha}\right)^{2 \alpha} \backslash \Delta_{\alpha}\right), x \geqslant 1-4 \alpha\right\} .
$$

Then, given constants $\epsilon \in(0,1 / 2)$ and $t \in \mathbb{R}$, we translate $\Phi_{\alpha}$ by $\epsilon \vec{e}_{2}+t \vec{e}_{3}$,

$$
\Phi_{\epsilon, \alpha}^{t}=\left\{(x, y+\epsilon, z+t):(x, y, z) \in \Phi_{\alpha}\right\} .
$$

Notice that definition of $\varphi_{\alpha}$ guarantees

$$
\Phi_{\epsilon, \alpha}^{t} \subset[1-4 \alpha, 1-\alpha] \times[-1+\epsilon,-1+\epsilon+2 \alpha] \times[-t-2 \alpha,+\infty) .
$$

Hence, there exists a constant $t_{\epsilon, \alpha} \in \mathbb{R}$ and a point $P_{0}=\left(x_{0}, y_{0}, z_{0}\right) \in \mathbb{R}^{3}$ satisfying

$$
\begin{equation*}
\Sigma \bigcap \Phi_{\epsilon, \alpha}^{t}=\varnothing \quad \text { for } t>t_{\epsilon, \alpha}, \quad P_{0} \in \Sigma \bigcap \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}} . \tag{4.20}
\end{equation*}
$$

We denote by $\Delta_{\epsilon, \alpha}$ the projection of $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ into the $x z$-plane, and consider $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ as the graph of a function $\varphi_{\epsilon, \alpha}: \Delta_{\epsilon, \alpha} \rightarrow \mathbb{R}$

$$
\Delta_{\epsilon, \alpha}=\left\{(x, z):(x, y, z) \in \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}\right\}, \quad \quad \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}=\left\{\left(x, \varphi_{\epsilon, \alpha}(x, z), z\right):(x, z) \in \Delta_{\epsilon, \alpha}\right\}
$$

Step 2 : Position of the contact point $P_{0}$. In this step, we will show that the contact point $P_{0}$ is contained in the front part $L_{1-4 \alpha}^{x}\left(\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}\right)$ of the boundary $\partial \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$.

First of all, the contact point $P_{0}$ can not be an interior point of $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$, because $\varphi_{\epsilon, \alpha}$ is a supersolution. Thus, $P_{0}$ is a point on the boundary $\partial \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ of $\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$. We observe that the boundary $\partial \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}$ can be decomposed into the top $\Gamma_{T}$, bottom $\Gamma_{B}$, and front $\Gamma_{F}$ boundary as following

$$
\begin{equation*}
\partial \Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}=\Gamma_{T} \bigcup \Gamma_{B} \bigcup \Gamma_{F}, \quad \Gamma_{T}=L_{-1+\epsilon+2 \alpha}^{y}\left(\Phi_{\epsilon, \alpha}^{t_{\epsilon}, \alpha}\right), \quad \Gamma_{B}=L_{-1+\epsilon}^{y}\left(\Phi_{\epsilon, \alpha}^{t_{\epsilon}, \alpha}\right), \quad \Gamma_{F}=L_{1-4 \alpha}^{x}\left(\Phi_{\epsilon, \alpha}^{t_{\epsilon, \alpha}}\right) . \tag{4.21}
\end{equation*}
$$

We denote by $P_{1}$ and $P_{2}$ the end point of the top $\Gamma_{T}$ and the bottom $\Gamma_{B}$ boundary, respectively

$$
\begin{align*}
& P_{1}=\left(x_{1}, y_{1}, z_{1}\right)=\Gamma_{T} \bigcap \Gamma_{F}=\left(1-4 \alpha,-1+\epsilon+2 \alpha, z_{1}\right),  \tag{4.22}\\
& P_{2}=\left(x_{2}, y_{2}, z_{2}\right)=\Gamma_{B} \bigcap \Gamma_{F}=\left(1-4 \alpha,-1+\epsilon, z_{2}\right) .
\end{align*}
$$

On the other hand (4.20) gives $\operatorname{cl}\left(\Delta_{\epsilon, \alpha}\right)=: \Delta_{\epsilon, \alpha} \subset \operatorname{Int}\left(\Omega_{y}\right)$ and $h(x, z) \leqslant \varphi_{\epsilon, \alpha}(x, z)$ on $\Delta_{\epsilon, \alpha}$. Also, we have $\left|D \varphi_{\epsilon, \alpha}\right|=+\infty$ on $\Gamma_{T}$. Hence, if $P_{0} \in\left(\Gamma_{T} \backslash\left\{P_{1}\right\}\right)$, then $|D h|=+\infty$ holds at $P_{0}$ by $h \leqslant \varphi_{\epsilon, \alpha}$, which contradicts to $\operatorname{cl}\left(\Delta_{\epsilon, \alpha}\right) \subset \operatorname{Int}\left(\Omega_{y}\right)$. Thus,

$$
P_{0} \notin\left(\Gamma_{T} \backslash\left\{P_{1}\right\}\right)
$$

Moreover, we have $\left|D \varphi_{\epsilon, \alpha}\right|=0$ on $\Gamma_{B}$. Thus, if $P_{0} \in\left(\Gamma_{B} \backslash\left\{P_{2}\right\}\right)$, then $|D h|=0$ holds at $P_{0}$. However, $\Sigma$ is a strictly convex complete surface, which means $|D h| \neq 0$. Therefore,

$$
P_{0} \notin\left(\Gamma_{B} \backslash\left\{P_{2}\right\}\right)
$$

Hence, by (4.21) and (4.22), $P_{0}$ is a point on the front boundary $\Gamma_{F}$

$$
\begin{equation*}
P_{0}=\left(x_{0}, y_{0}, z_{0}\right)=\left(1-4 \alpha, y_{0}, z_{0}\right) \in \Gamma_{F} . \tag{4.23}
\end{equation*}
$$

Step 3 : Distance between $P_{0}$ and $P_{2}$. In this step, we will estimate $z_{2}-z_{0}$ in terms of $\alpha$.
We recall that the Grim reaper curve $\partial \Delta_{\alpha}$ is the graph of the function $f_{\alpha}(x)$ defined by (4.18) and $\Delta_{\epsilon, \alpha}$ is a subset of $\bar{\Delta}_{\epsilon, \alpha}=: \Delta_{\alpha}+t_{\epsilon, \alpha} e_{z}$. Hence, $\partial \bar{\Delta}_{\epsilon, \alpha}$ is the graph of the function $f_{\epsilon, \alpha}$ defined by

$$
\begin{equation*}
f_{\epsilon, \alpha}(x)=t_{\epsilon, \alpha}+f_{\alpha}(x)=t_{\epsilon, \alpha}-\frac{2}{\pi \alpha} \log \cos \left(\frac{\alpha \pi}{2}\left(x-1+3 \alpha+\frac{1}{\alpha}\right)\right) . \tag{4.24}
\end{equation*}
$$

By definition of $\varphi_{\alpha}$, there exists a unique point $\left(\bar{x}_{0}, \bar{z}_{0}\right) \in \partial \Delta_{\epsilon, \alpha}$ such that

$$
\begin{equation*}
d\left(\left(x_{0}, z_{0}\right), \bar{\Delta}_{\epsilon, \alpha}\right)=d\left(\left(x_{0}, z_{0}\right),\left(\bar{x}_{0}, \bar{z}_{0}\right)\right) \leqslant 2 \alpha \tag{4.25}
\end{equation*}
$$

We know $x_{0}=x_{2}=1-4 \alpha$ by (4.22) and (4.23). Hence, for all $x \in\left[\bar{x}_{0}, x_{2}\right]$, we can derive from (4.24) the following inequality

$$
\begin{equation*}
f_{\epsilon, \alpha}^{\prime}\left(\bar{x}_{0}\right) \leqslant f_{\epsilon, \alpha}^{\prime}(x) \leqslant f_{\epsilon, \alpha}^{\prime}\left(x_{2}\right)=\tan \left(\frac{\pi \alpha}{2}\left(x_{0}-1+3 \alpha+\frac{1}{\alpha}\right)\right)=\cot \frac{\pi \alpha^{2}}{2} . \tag{4.26}
\end{equation*}
$$

Therefore, combining (4.25) and (4.26) yields

$$
\begin{equation*}
z_{2}-z_{0} \leqslant\left(\bar{z}_{0}-z_{0}\right)+\left(z_{2}-\bar{z}_{0}\right) \leqslant 2 \alpha+\left(z_{6}-\bar{z}_{0}\right) \leqslant 2 \alpha+\int_{\bar{x}_{0}}^{x_{2}} f_{\epsilon, \alpha}^{\prime}(x) d x \leqslant 2 \alpha+2 \alpha \cot \frac{\pi \alpha^{2}}{2} \tag{4.27}
\end{equation*}
$$



Figure 4. Level sets of the solution $h$
Step 4 : Partial derivative $\partial_{x} u$ bound at the contact point $P_{0}$. We can consider the level set $L_{y_{0}}\left(\varphi_{\epsilon, \alpha}\right)$ as the graph of a convex function $f_{\epsilon, \alpha}^{0}:[1-4 \alpha, 1-\alpha) \rightarrow \mathbb{R}$, namely $L_{y_{0}}\left(\varphi_{\epsilon, \alpha}\right)=$ $\left\{\left(x, f_{\epsilon, \alpha}^{0}(x)\right): x \in[1-4 \alpha, 1-\alpha)\right\}$. Then, by the definitions of $\varphi_{\alpha}$ and $\left(\bar{x}_{0}, \bar{z}_{0}\right)$, we have $\left(f_{\epsilon, \alpha}^{0}\right)^{\prime}\left(x_{0}\right)=f_{\epsilon, \alpha}^{\prime}\left(\bar{x}_{0}\right)$. Thus, 4.26 yields the bound

$$
\left(f_{\epsilon, \alpha}^{0}\right)^{\prime}\left(x_{0}\right) \leqslant \cot \left(\pi \alpha^{2} / 2\right) .
$$

On the other hand, 4.20) implies $L_{y_{0}}\left(\varphi_{\epsilon, \alpha}\right)<L_{y_{0}}(h)$, namely $u\left(x, y_{0}\right) \leqslant f_{\epsilon, \alpha}^{0}(x)$ holds for all $x \in[1-4 \alpha, 1-\alpha)$. Therefore,

$$
\begin{equation*}
\partial_{x} u\left(x_{0}, y_{0}\right) \leqslant\left(f_{\epsilon, \alpha}^{0}\right)^{\prime}\left(x_{0}\right) \leqslant \cot \left(\pi \alpha^{2} / 2\right) . \tag{4.28}
\end{equation*}
$$

Step 5 : Partial derivative $\partial_{x} u$ bound at the given point. We define the points $P_{3}, P_{4}, P_{5}$ on $\Sigma$ by

$$
\begin{aligned}
& P_{3}=\left(x_{3}, y_{3}, z_{3}\right)=:\left(1-5 \alpha, y_{0}, u\left(1-5 \alpha, y_{0}\right)\right), \\
& P_{4}=\left(x_{4}, y_{4}, z_{4}\right)=:(1-5 \alpha,-1+\epsilon, u(1-5 \alpha,-1+\epsilon)), \\
& P_{5}=\left(x_{5}, y_{5}, z_{5}\right)=:(1-4 \alpha,-1+\epsilon, u(1-4 \alpha,-1+\epsilon)) .
\end{aligned}
$$

Since we know $P_{0}, P_{3} \in L_{y_{0}}^{y}(\Sigma)$, the inequality (4.28) and the convexity of $u$ give

$$
z_{0}-z_{3}=\int_{x_{3}}^{x_{0}} \partial_{x} u\left(x, y_{0}\right) d x \leqslant \int_{x_{3}}^{x_{0}} \partial_{x} u\left(x_{0}, y_{0}\right) d x=\alpha\left(\partial_{x} u\left(x_{0}, y_{0}\right)\right) \leqslant \alpha \cot \frac{\pi \alpha^{2}}{2} .
$$

By adding (4.27) and the inequality above, we obtain

$$
\begin{equation*}
z_{2}-z_{3} \leqslant 2 \alpha+3 \alpha \cot \left(\pi \alpha^{2} / 2\right) \tag{4.29}
\end{equation*}
$$

On the other hand, 4.20) implies that $L_{-1+\epsilon}\left(\varphi_{\epsilon, \alpha}\right)<L_{-1+\epsilon}(h)$. Therefore, $x_{2}=x_{5}=1-4 \alpha$, $\left(x_{2}, z_{2}\right) \in L_{-1+\epsilon}\left(\varphi_{\epsilon, \alpha}\right)$, and $\left(x_{5}, z_{5}\right) \in L_{-1+\epsilon}(h)$ guarantee

$$
z_{2}=\varphi_{\epsilon, \alpha}\left(x_{2}\right) \geqslant u\left(x_{2},-1+\epsilon\right)=u\left(x_{5},-1+\epsilon\right)=z_{5} .
$$

Also, the convexity and symmetry of $\Sigma$ give that

$$
z_{3} \leqslant z_{4}
$$

Thus, subtracting the inequalities above yields $z_{5}-z_{4} \leqslant z_{2}-z_{3}$. Applying 4.29), we have

$$
z_{5}-z_{4} \leqslant 2 \alpha+3 \alpha \cot \left(\pi \alpha^{2} / 2\right)
$$

Hence, the desired result follows by the following computation

$$
z_{5}-z_{4}=\int_{x_{4}}^{x_{5}} \partial_{x} u(x,-1+\epsilon) d x \geqslant \int_{x_{4}}^{x_{5}} \partial_{x} u\left(x_{4},-1+\epsilon\right) d x=\alpha\left(\partial_{x} u(1-5 \alpha,-1+\epsilon)\right) .
$$

## 3 Distance from the tip to flat sides

Let $\Sigma$ be the translating solution to the Gauss curvature flow over the square $\Omega$ as in Theorem 4.2. In this final section we will show that this solution has flat sides, as stated in Theorem 4.12 below. To this end, we will study the distance from the tip of a solution $\Sigma$ to each point on the free boundary, that is the boundary of the flat sides. To estimate this distance one needs to establish a gradient bound for solutions to the equation (4.1) at a certain point near the flat sides. Since the gradient bound depends on the global structure of $\Omega$, we will establish an integral estimate by deriving a separation of variables structure from (4.1) as in the proof of the following Lemma.

Lemma 4.11 (Gradient bound). Let $\Omega$ satisfy the conditions in Theorem 4.10 and let $u$ be a solution of (4.1) on $\Omega$. Assume that $[a, b] \times[-1,-1+\sigma] \subset \Omega$, for some constants $a, b, \sigma \in(0,1)$. Then, there exists a point $x_{0} \in[a, b]$ satisfying

$$
-\partial_{y} u\left(x_{0},-1+\sigma\right) \leqslant \sqrt{\frac{2 M}{\pi \sigma(b-a)}},
$$

where $M=\sup _{y \in(0, \sigma)} \partial_{x} u(b,-1+y)$.

Proof. Since the domain $\omega$ satisfies the assumptions of Theorem4.10, we have $\mathcal{A} \leqslant 64 / 9$. Hence, the following inequality holds

$$
\frac{u_{y y} u_{x x}}{\left(1+u_{y}^{2}\right)^{\frac{3}{2}}} \geqslant \frac{u_{y y} u_{x x}-u_{x y}^{2}}{\left(1+u_{x}^{2}+u_{y}^{2}\right)^{\frac{3}{2}}}=\frac{\operatorname{det} D^{2} u}{\left(1+|D u|^{2}\right)^{\frac{3}{2}}}=\frac{2 \pi}{\mathcal{A}} \geqslant \frac{\pi}{4} .
$$

which combined with Holder inequality yields

$$
\left(\int_{a}^{b} \frac{u_{y y}}{\left(1+u_{y}^{2}\right)^{\frac{3}{2}}} d x\right)\left(\int_{a}^{b} u_{x x} d x\right) \geqslant\left(\int_{a}^{b}(\pi / 4)^{\frac{1}{2}} d x\right)^{2}=\frac{\pi}{4}(a-b)^{2} .
$$

Since $a \geqslant 0$ and Proposition 4.7 guarantee that $u_{x}(\cdot, a) \geqslant 0$, we have

$$
\int_{a}^{b} u_{x x} d x=u_{x}(y, b)-u_{x}(y, a) \leqslant u_{x}(y, b) \leqslant M
$$

Hence,

$$
\frac{1}{b-a} \int_{a}^{b} \int_{-1}^{-1+\sigma} \frac{u_{y y}}{\left(1+u_{y}^{2}\right)^{\frac{3}{2}}} d y d x \geqslant \frac{\pi \sigma}{4 M}(b-a)
$$

Therefore, there exists a constant $x_{0} \in[a, b]$ satisfying

$$
\int_{-1}^{-1+\sigma} \frac{u_{y y}}{\left(1+u_{y}^{2}\right)^{\frac{3}{2}}}\left(x_{0}, \cdot\right) d y \geqslant \frac{\pi \sigma}{4 M}(b-a) .
$$

Moreover, by $\left|u_{y}\right| \leqslant\left(1+u_{y}^{2}\right)^{\frac{1}{2}}$, we have

$$
\int_{-1}^{-1+\sigma} \frac{u_{y y}}{\left(1+u_{y}^{2}\right)^{\frac{3}{2}}}\left(x_{0}, \cdot\right) d y=\left.\frac{u_{y}}{\left(1+u_{y}^{2}\right)^{\frac{1}{2}}}\left(x_{0}, \cdot\right)\right|_{-1} ^{-1+\sigma} \leqslant \frac{u_{y}}{\left(1+u_{y}^{2}\right)^{\frac{1}{2}}}\left(x_{0},-1+\sigma\right)+1 .
$$

Hence, at $\left(x_{0},-1+\sigma\right)$, the following holds

$$
\frac{\left|u_{y}\right|}{\left(1+u_{y}^{2}\right)^{\frac{1}{2}}}\left(x_{0},-1+\sigma\right)=\frac{-u_{y}}{\left(1+u_{y}^{2}\right)^{\frac{1}{2}}}\left(x_{0},-1+\sigma\right) \leqslant 1-\frac{\pi \sigma}{4 M}(b-a)
$$

implying the bound

$$
1+u_{y}^{-2}\left(x_{0},-1+\sigma\right) \geqslant\left(1-\frac{\pi \sigma}{4 M}(b-a)\right)^{-2} \geqslant\left(1+\frac{\pi \sigma}{4 M}(b-a)\right)^{2} \geqslant 1+\frac{\pi \sigma}{2 M}(b-a) .
$$

Solving this last inequality for $-u_{y}\left(x_{0},-1+\sigma\right) \geqslant 0$ leads to the desired result.
The distance between the tip of the translating solution $\Sigma$ over the square and its flat sides is estimated in the following result.

Theorem 4.12 (Distance between the tip and flat sides). Let $\Omega$ satisfy the conditions in Theorem 4.10 and let $u$ be a solution of (4.1). Given $\alpha \in(0,1 / 6)$, there exists a constant $C>0$ satisfying

$$
u(1-6 \alpha,-1)-u(0,0) \leqslant 6\left(1+\frac{1}{\alpha^{2}}\right)
$$



Figure 5. Converging points on domain $\Omega$

Proof. We begin by setting $a_{0}=1-6 \alpha, b_{0}=1-5 \alpha, \sigma_{0}=\frac{1}{2}$, and $M=\sup _{y \in(0,1 / 2)} \partial_{x} u\left(b_{0},-1+\right.$ $y)$. Then, by Lemma 4.11, there exists a point $x_{0} \in\left[a_{0}, b_{0}\right]$ satisfying

$$
\begin{equation*}
-\partial_{y} u\left(x_{0},-1+1 / 2\right) \leqslant 2(M / \pi \alpha)^{\frac{1}{2}} . \tag{4.30}
\end{equation*}
$$

We choose an interval $\left[a_{1}, b_{1}\right]$ satisfying $x_{0} \in\left[a_{1}, b_{1}\right] \subset\left[a_{0}, b_{0}\right]$ and $b_{1}-a_{1}=2^{-1 / 3} \alpha$. Then, for $\sigma_{1}=2^{-2}$, we have $\sup _{y \in\left(0, \sigma_{1}\right)} \partial_{x} u\left(b_{1},-1+y\right) \leqslant M$. Hence, Lemma 4.11 gives a point $x_{1} \in\left[a_{1}, b_{1}\right]$ satisfying

$$
-\partial_{y} u\left(x_{1},-1+1 / 2^{2}\right) \leqslant 2^{1+\frac{2}{3}}(M / \pi \alpha)^{\frac{1}{2}} .
$$

By setting $\sigma_{n}=2^{-1-n}$, we can inductively choose intervals [ $a_{n}, b_{n}$ ] satisfying $x_{n-1} \in\left[a_{n}, b_{n}\right] \subset$ [ $\left.a_{n-1}, b_{n-1}\right]$ and $b_{n}-a_{n}=2^{-n / 3} \alpha$ so that we obtain a point $x_{n} \in\left[a_{n}, b_{n}\right]$ satisfying

$$
-\partial_{y} u\left(x_{n},-1+1 / 2^{n+1}\right) \leqslant 2^{1+\frac{2 n}{3}}(M / \pi \alpha)^{\frac{1}{2}} .
$$

Then, integrating along $y$ yields

$$
\left|u\left(x_{n},-1+\frac{1}{2^{n+1}}\right)-u\left(x_{n},-1+\frac{1}{2^{n}}\right)\right| \leqslant \int_{-1+1 / 2^{n+1}}^{-1+1 / 2^{n}}-\partial_{y} u\left(x_{n}, y\right) d y \leqslant 2^{-n / 3}(M / \pi \alpha)^{\frac{1}{2}}
$$

On the other hand, $x_{n-1}, x_{n} \in\left[a_{n}, b_{n}\right]$ and $a_{n}-b_{n}=2^{-n / 3} \alpha$ imply

$$
\left|u\left(x_{n},-1+\frac{1}{2^{n}}\right)-u\left(x_{n-1},-1+\frac{1}{2^{n}}\right)\right|=\left|\int_{x_{n-1}}^{x_{n}} \partial_{x} u\left(x,-1+\frac{1}{2^{n}}\right) d x\right| \leqslant 2^{-\frac{n}{3}} M \alpha
$$

Therefore,

$$
\left|u\left(x_{n},-1+\frac{1}{2^{n+1}}\right)-u\left(x_{n-1},-1+\frac{1}{2^{n}}\right)\right| \leqslant 2^{-\frac{n}{3}}\left((M / \pi \alpha)^{\frac{1}{2}}+M \alpha\right) .
$$

By definition of $x_{n}$, the sequence $\left\{x_{n}\right\}_{n \in \mathbb{N}}$ converges to a point $\bar{x} \in\left[a_{0}, b_{0}\right]$. Hence, we can sum up the inequality above for all $n \in \mathbb{N}$ so that we have

$$
\begin{equation*}
u(\bar{x},-1)-u\left(x_{0},-\frac{1}{2}\right) \leqslant \sum_{n=1}^{\infty} 2^{-\frac{n}{3}}\left((M / \pi \alpha)^{\frac{1}{2}}+M \alpha\right) \leqslant 4\left((M / \pi \alpha)^{\frac{1}{2}}+M \alpha\right) \tag{4.31}
\end{equation*}
$$

Next, we consider the linear function

$$
f(x, y)=\left(\partial_{x} u\left(x_{0},-1 / 2\right)\right)\left(x-x_{0}\right)+\left(\partial_{y} u\left(x_{0},-1 / 2\right)\right)(y+1 / 2)+u\left(x_{0},-1 / 2\right)
$$

whose graph is the tangent hyperplane of $\Sigma$ at $\left(x_{0},-1 / 2, u\left(x_{0},-1 / 2\right)\right)$. Then, the convexity of $\Sigma$ gives $f(0,0) \leqslant u(0,0)$. Hence, (4.30) and definition of $M$ show

$$
u\left(x_{0},-\frac{1}{2}\right)-u(0,0) \leqslant f\left(x_{0},-\frac{1}{2}\right)-f(0,0) \leqslant(M / \pi \alpha)^{\frac{1}{2}}+M
$$

Thus, Theorem 4.10, (4.31), and the inequality give

$$
u(\bar{x},-1)-u(0,0) \leqslant 5\left(\frac{M}{\pi \alpha}\right)^{\frac{1}{2}}+2 M \leqslant 5\left(\frac{2}{\pi \alpha}+\frac{3}{\pi \alpha} \cot \frac{\pi \alpha^{2}}{2}\right)^{\frac{1}{2}}+2\left(2+3 \cot \frac{\pi \alpha^{2}}{2}\right)
$$

Applying $1 / \alpha \geqslant 6$ and $\cot \left(\pi \alpha^{2} / 2\right) \leqslant 2 /\left(\pi \alpha^{2}\right)$, we have

$$
u(\bar{x},-1)-u(0,0) \leqslant 5\left(\frac{2}{\pi \alpha}+\frac{6}{\pi^{2} \alpha^{3}}\right)^{\frac{1}{2}}+\frac{12}{\pi \alpha^{2}}+4 \leqslant 6\left(1+\frac{1}{\alpha^{2}}\right) .
$$

Finally, combining Proposition 4.7, the convexity of $u$, and $\bar{x} \geqslant b_{0}=1-6 \alpha \geqslant 0$ yields $u(\bar{x},-1) \geqslant u(1-6 \alpha,-1)$, which leads to the desired result.

We will now give the proof of the main Theorem4.2. This redily follows from the following result.

Theorem 4.13 (Existence of flat sides). Let $\Omega=(-1,1) \times(-1,1)$ and $u$ be a solution of (4.1) on $\Omega$. Then, there exists a function $\bar{u}:(\partial \Omega \backslash V) \rightarrow \mathbb{R}$ satisfying

$$
\bar{u}\left(x_{0}\right)=\lim _{x \rightarrow x_{0}} u(x) .
$$

Also, the corresponding complete solution $\Sigma$ of (4.2) is a convex surface of class $C_{l o c}^{1,1}$, and $u$ satisfies

$$
\lim _{x \rightarrow V} u(x)=+\infty .
$$

Proof. Let $\left\{\Omega_{n}\right\}_{n \in \mathbb{N}}$ be a sequence of sets satisfying the conditions in Theorem 4.10, and let $\left\{u_{n}\right\}_{n \in \mathbb{N}}$ and $\left\{\Sigma_{n}\right\}_{n \in \mathbb{N}}$ be the sequence of corresponding solutions of (4.1) with $u_{n}(0,0)=0$ and their graphs, respectively. We denote the convex hull of $\Sigma_{n}$ by $E_{n}=\{t X+(1-t) Y: X, Y \in$ $\left.\Sigma_{n}, t \in[0,1]\right\}$, and define a convex body $E$ by

$$
E=\bigcap_{n \in \mathbb{N}} E_{n} .
$$

Observe that $E$ is not an empty set, because Theorem 4.12 gives that $\left(1-6 \alpha,-1,6\left(1+\alpha^{-2}\right)\right) \in E_{n}$ which means that $\left(1-6 \alpha,-1,6\left(1+\alpha^{-2}\right)\right) \in E$. Now, we denote by $\Sigma$ the boundary of $E$. Then, $\Sigma$ is naturally a convex complete and non-compact surface, since $\{(0,0, t): t \geqslant 0\} \subset E_{n}$ implies $\{(0,0, t): t \geqslant 0\} \subset E$.

We define $u: \Omega \rightarrow \mathbb{R}$ by $(x, y, u(x, y)) \in \Sigma$. Then, the symmetry of $\Omega_{n}$, the convexity of $E$, and $\left(1-6 \alpha,-1,6\left(1+\alpha^{-2}\right)\right) \in E$ guarantees that there exists a function $\bar{u}:(\partial \Omega \backslash V) \rightarrow \mathbb{R}$ satisfying

$$
\bar{u}\left(x_{0}\right)=\lim _{x \rightarrow x_{0}} u(x) .
$$

Moreover, Theorem 4.4 shows the local $C^{1,1}$ regularity of $\Sigma$. Now, we assume that there exists a point $\left(1,1, t_{0}\right)$ in $E$, then the mean curvature of $\Sigma$ attains $+\infty$ at $(1,1, t)$ for $t>t_{0}$. This contradicts to the local $C^{1,1}$ regularity of $\Sigma$. Hence, we have $\Sigma \bigcap(V \times \mathbb{R})=\varnothing$, namely $u$ satisfies

$$
\lim _{x \rightarrow V} u(x)=+\infty .
$$

## Chapter 5

## Asymptotic Behavior of closed solutions

In this chapter, we show the following Theorem proven in [9].
Theorem 5.1 (Choi-Daskalopoulos). Given $\alpha \in\left(\frac{1}{n}, 1+\frac{1}{n}\right)$, the unit $n$-sphere is the unique closed strictly convex smooth solution to $\left({ }^{\left.* *^{\alpha}\right)}\right.$.

Notice that jointly with Brendle, the result was extend to $\alpha \geqslant \frac{1}{n+2}$ in [5].
Theorem 5.2 (Brendle-Choi-Daskalopoulos). Given $\alpha>\frac{1}{n+2}$, the unit $n$-sphere is the unique closed strictly convex smooth solution to $\left(\sqrt{\left.* *^{\alpha}\right)}\right.$. If $\alpha=\frac{1}{n+2}$, closed strictly convex smooth solutions to $\left(* *^{* *^{\alpha}}\right)$ are ellipsoids.

Theorem 5.2 combined with the results in [4, 19] imply the convergence of the $\alpha$-Gauss curvature flow to the round sphere (or ellipsoids), which in particular proves the higher dimensional Firey's conjecture.

Theorem 5.3. Let $\Sigma_{t}$ be a strictly convex, closed and smooth solution to the $\alpha$-Gauss curvature flow with $\alpha \in \frac{1}{n+2}, n \geqslant 2$. Then, there exists a finite time $T$ at which the solution $\Sigma_{t}$ converges after rescaling to the round sphere. If $\alpha=\frac{1}{n+2}$, the solution converges after rescaling to an ellipsoid.

Remark 5.4 (Pogorelov estimate on powers of a matrix). Pogorelov type estimates in this context have been frequently applied in the past by using $b^{11}$, the first entry of a matrix $A^{-1}:=$ $\left\{b^{i j}\right\}$. However, if one applies the Pogorelov estimate for $b^{11} K^{\alpha}-\frac{n \alpha-1}{2 n \alpha}|F|^{2}$, one can obtain the result of Theorem 5.5 only for $\alpha \in\left(\frac{1}{n}, \frac{1}{2}\right]$. In this work, by using instead $\left(b^{1 i} g_{i j} b^{j 1}\right)^{\frac{1}{2}}$, the root of the first entry of the square $A^{-2}$ of the matrix $A^{-1}$, we are able to extend the result of Theorem 5.5 to the range of exponents $\alpha \in\left(\frac{1}{n}, 1+\frac{1}{n}\right)$, which includes the classical case of the Gauss curvature flow $\alpha=1$.

In [5], a Pogorelov type estimate was developed in viscosity sense which extends the result of Theorem 5.5 for all $\alpha>\frac{1}{n}$.

## 1 Pogorelov type computation

We consider the function $W: M^{n} \rightarrow \mathbb{R}$ given by

$$
W(p):=\left(K^{\alpha} \lambda_{\min }^{-1}-\frac{n \alpha-1}{2 n \alpha}|F|^{2}\right)(p) .
$$

We will employ in this section a Pogorelov type computation to show that the maximum point of $W(p)$ is an umbilical point. We begin with the following standard observation which we include here for the reader's convenience.

We will now show that one of the Pogorelov type expressions of the function $W$ plays a role as a subsolution of $\left(\sqrt{* *^{\alpha}}\right)$ at a given maximum point, to imply that the maximum point of $W(p)$ is an umbilical point.

Theorem 5.5 (Pogorelov type computation). Let $\Sigma$ be a strictly convex smooth closed solution of $\left({* *^{\alpha}}^{\alpha}\right)$ for an exponent $\alpha \in\left(\frac{1}{n}, 1+\frac{1}{n}\right)$. Assume that $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ is a smooth immersion such that $F\left(M^{n}\right)=\Sigma$, and the continuous function $W(p)$ attains its maximum at a point $p_{0}$. Then, $F\left(p_{0}\right)$ is an umbilical point and $\nabla|F|^{2}\left(p_{0}\right)=0$ holds.

Proof. We begin by choosing a coordinate chart $(U, \varphi)$ of $p_{0} \in \varphi(U) \subset M^{n}$ such that the covariant derivatives $\left\{\nabla_{i} F\left(p_{0}\right):=\partial_{i}(F \circ \varphi)\left(\varphi^{-1}\left(p_{0}\right)\right)\right\}_{i=1, \cdots, n}$ form an orthonormal basis of $T \Sigma_{F\left(p_{0}\right)}$ satisfying

$$
g_{i j}\left(p_{0}\right)=\delta_{i j}, \quad h_{i j}\left(p_{0}\right)=\delta_{i j} \lambda_{i}\left(p_{0}\right), \quad \lambda_{1}\left(p_{0}\right)=\lambda_{\min }\left(p_{0}\right),
$$

which guarantees $b^{11}\left(p_{0}\right)=\lambda_{\min }^{-1}\left(p_{0}\right)$ and $g^{11}\left(p_{0}\right)=1$. Next, we define the function $\bar{W}: \varphi(U) \rightarrow$ $\mathbb{R}$ by

$$
\bar{W}(p):=K^{\alpha}\left(\frac{b^{1 i} g_{i j} b^{j 1}}{g^{11}}\right)^{\frac{1}{2}}(p)-\frac{n \alpha-1}{2 n \alpha}|F|^{2}(p) .
$$

Then, by Proposition 2.5 we have

$$
\bar{W}(p) \leqslant W(p) \leqslant W\left(p_{0}\right)=\bar{W}\left(p_{0}\right)
$$

which means that $\bar{W}$ attains its maximum at $p_{0}$.
We will now calculate $\mathcal{L} \bar{W}:=\alpha K^{\alpha} b^{i j} \nabla_{i} \nabla_{j} \bar{W}$ at the point $p_{0}$. First we derive the following equation from (2.18)

$$
\begin{aligned}
\mathcal{L}\left(b_{p}^{1} b^{p 1}\right)= & 2 \alpha K^{\alpha} b^{i j} \nabla_{i} b^{p 1} \nabla_{j} b_{p}^{1}+2 K^{-\alpha} b_{p}^{1} b^{p r} b^{1 s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}+2 \alpha K^{\alpha} b_{p}^{1} b^{p r} b^{1 s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m} \\
& +\left\langle F, \nabla\left(b_{p}^{1} b^{p 1}\right)\right\rangle-2 b_{p}^{1} b^{p 1}-2(n \alpha-1) K^{\alpha} b^{11}+2 \alpha K^{\alpha} H b_{p}^{1} b^{p 1} .
\end{aligned}
$$

Thus, we obtain

$$
\begin{aligned}
\mathcal{L}\left(\frac{b_{p}^{1} b^{p 1}}{g^{11}}\right)^{\frac{1}{2}}= & -\frac{\alpha K^{\alpha} b^{i j} \nabla_{i}\left(b_{p}^{1} b^{1 p}\right) \nabla_{j}\left(b_{q}^{1} b^{1 q}\right)}{4\left(b_{r}^{1} b^{r 1}\right)^{\frac{3}{2}}\left(g^{11}\right)^{\frac{1}{2}}}+\frac{\alpha K^{\alpha} b^{i j} \nabla_{i} b^{p 1} \nabla_{j} b_{p}^{1}}{\left(b_{q}^{1} b^{q 1} g^{11}\right)^{\frac{1}{2}}} \\
& +\frac{b_{p}^{1} b^{p r} b^{1 s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}}{K^{\alpha}\left(b_{q}^{1} b^{q 1} g^{11}\right)^{\frac{1}{2}}}+\frac{\alpha K^{\alpha} b_{p}^{1} b^{p r} b^{1 s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m}}{\left(b_{q}^{1} b^{q 1} g^{11}\right)^{\frac{1}{2}}} \\
& +\left\langle F, \nabla\left(b_{p}^{1} b^{p 1} / g^{11}\right)^{\frac{1}{2}}\right\rangle-\left(\frac{b_{p}^{1} b^{p 1}}{g^{11}}\right)^{\frac{1}{2}}-\frac{(n \alpha-1) K^{\alpha} b^{11}}{\left(b_{p}^{1} b^{p 1} g^{11}\right)^{\frac{1}{2}}}+\alpha K^{\alpha} H\left(\frac{b_{p}^{1} b^{p 1}}{g^{11}}\right)^{\frac{1}{2}} .
\end{aligned}
$$

Combining this with (2.17) yields

$$
\begin{align*}
\mathcal{L} \bar{W}= & -\frac{n \alpha-1}{2 n \alpha} \mathcal{L}|F|^{2}+2\left\langle\nabla K^{\alpha}, \nabla\left(\frac{b_{p}^{1} b^{p 1}}{g^{11}}\right)^{\frac{1}{2}}\right\rangle_{\mathcal{L}}-\frac{\alpha K^{2 \alpha} b^{i j} b_{p}^{1} b_{q}^{1} \nabla_{i} b^{1 p} \nabla_{j} b^{1 q}}{\left(b_{r}^{1} b^{r 1}\right)^{\frac{3}{2}}\left(g^{11}\right)^{\frac{1}{2}}}  \tag{5.1}\\
& +\frac{\alpha K^{2 \alpha} b^{i j} \nabla_{i} b^{p 1} \nabla_{j} b_{p}^{1}}{\left(b_{q}^{1} b^{q 1} g^{11}\right)^{\frac{1}{2}}}+\frac{b_{p}^{1} b^{p r} b^{1 s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}}{\left(b_{q}^{1} b^{q 1} g^{11}\right)^{\frac{1}{2}}}+\frac{\alpha K^{2 \alpha} b_{p}^{1} b^{p r} b^{1 s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m}}{\left(b_{q}^{1} b^{1} g^{11}\right)^{\frac{1}{2}}} \\
& +\left\langle F, \nabla\left(K^{\alpha}\left(b_{p}^{1} b^{p 1} / g^{11}\right)^{\frac{1}{2}}\right)\right\rangle+(n \alpha-1) K^{\alpha}\left(\frac{b_{p}^{1} b^{p 1}}{g^{11}}\right)^{\frac{1}{2}}-\frac{(n \alpha-1) K^{2 \alpha} b^{11}}{\left(b_{p}^{1} b^{p 1} g^{11}\right)^{\frac{1}{2}}} .
\end{align*}
$$

Observe that

$$
2\left\langle\nabla K^{\alpha}, \nabla\left(\frac{b_{p}^{1} b^{p 1}}{g^{11}}\right)^{\frac{1}{2}}\right\rangle_{\mathcal{L}}=2 \alpha K^{\alpha} b^{i j}\left(g^{11}\right)^{-\frac{1}{2}}\left(b_{q}^{1} b^{q 1}\right)^{-\frac{1}{2}} b_{p}^{1} \nabla_{i} K^{\alpha} \nabla_{j} b^{p 1},
$$

and

$$
\nabla\left(K^{\alpha}\left(b_{p}^{1} b^{p 1} / g^{11}\right)^{\frac{1}{2}}\right)=\nabla \bar{W}+\frac{n \alpha-1}{2 n \alpha} \nabla|F|^{2}
$$

Hence, applying the equations above, (2.15) and $\nabla \bar{w}\left(p_{0}\right)=0$ to (5.1) yields that the following holds at the maximum point $p_{0}$

$$
\begin{align*}
0 \geqslant & 2 \alpha K^{\alpha} \sum_{i=1}^{n} b^{i i} \nabla_{i} K^{\alpha} \nabla_{i} b^{11}-\alpha K^{2 \alpha} \sum_{i=1}^{n} b^{i i} h_{11}\left|\nabla_{i} b^{11}\right|^{2}+\alpha K^{2 \alpha} \sum_{j, p} b^{j j} h_{11}\left|\nabla_{j} b^{p 1}\right|^{2}+\left|b^{11} \nabla_{1} K^{\alpha}\right|^{2}  \tag{5.2}\\
& \left.+\alpha K^{2 \alpha}\left(b^{11}\right)^{2} \sum_{i, j} b^{i i} b^{j j}\left|\nabla_{1} h_{i j}\right|^{2}+\left.\frac{n \alpha-1}{2 n \alpha}\langle F, \nabla| F\right|^{2}\right\rangle+(n \alpha-1) K^{\alpha}\left(b^{11}-\frac{1}{n} \sum_{i=1}^{n} b^{i i}\right)
\end{align*}
$$

By (2.14), the second and third terms on the right hand side of the inequality above (5.2) satisfy

$$
\begin{aligned}
& -\sum_{i=1}^{n} b^{i i} h_{11}\left|\nabla_{i} b^{11}\right|^{2}+\sum_{j, p} b^{j j} h_{11}\left|\nabla_{j} b^{p 1}\right|^{2}=-\sum_{i=1}^{n} b^{i i}\left(b^{11}\right)^{3}\left|\nabla_{i} h_{11}\right|^{2}+\sum_{j, p} b^{j j} b^{11}\left(b^{p p}\right)^{2}\left|\nabla_{j} h_{p 1}\right|^{2} \\
& =\sum_{j=1}^{n} \sum_{p \neq 1} b^{j j} b^{11}\left(b^{p p}\right)^{2}\left|\nabla_{j} h_{p 1}\right|^{2} \geqslant \sum_{p \neq 1}\left(b^{11} b^{p p}\right)^{2}\left|\nabla_{p} h_{11}\right|^{2}=\sum_{p \neq 1}\left(b^{p p} h_{11}\right)^{2}\left|\nabla_{p} b^{11}\right|^{2}
\end{aligned}
$$

Also, by (2.14) the fifth term on the right hand side of (5.2) satisfies

$$
\left(b^{11}\right)^{2} \sum_{i, j} b^{i i} b^{j j}\left|\nabla_{1} h_{i j}\right|^{2} \geqslant\left(b^{11}\right)^{4}\left|\nabla_{1} h_{11}\right|^{2}+2 \sum_{i \neq 1}\left(b^{11}\right)^{3} b^{i i}\left|\nabla_{i} h_{11}\right|^{2}=\left|\nabla_{1} b^{11}\right|^{2}+2 \sum_{i \neq 1} b^{i i} h_{11}\left|\nabla_{i} b^{11}\right|^{2} .
$$

Furthermore, we have

$$
\alpha K^{2 \alpha}\left|\nabla_{1} b^{11}\right|^{2}+2 \alpha K^{\alpha} b^{11} \nabla_{1} K^{\alpha} \nabla_{1} b^{11} \geqslant-\alpha\left|b^{11} \nabla_{1} K^{\alpha}\right|^{2} .
$$

Hence, by applying the inequalities above, we can reduce (5.2) to

$$
\begin{align*}
0 \geqslant & 2 \alpha \sum_{i \neq 1} b^{i i} \nabla_{i} K^{\alpha}\left(K^{\alpha} \nabla_{i} b^{11}\right)+\alpha \sum_{p \neq 1}\left(b^{p p} h_{11}\right)^{2}\left|K^{\alpha} \nabla_{p} b^{11}\right|^{2}+2 \alpha \sum_{i \neq 1} b^{i i} h_{11}\left|K^{\alpha} \nabla_{i} b^{11}\right|^{2}  \tag{5.3}\\
& \left.+(1-\alpha)\left|b^{11} \nabla_{1} K^{\alpha}\right|^{2}+\left.\frac{n \alpha-1}{2 n \alpha}\langle F, \nabla| F\right|^{2}\right\rangle+(n \alpha-1) K^{\alpha}\left(b^{11}-\frac{1}{n} \sum_{i=1}^{n} b^{i i}\right) .
\end{align*}
$$

We now employ (2.16) to obtain the following at the point $p_{0}$

$$
\begin{equation*}
b^{i i} \nabla_{i} K^{\alpha}=b^{i i} h_{i i}\left\langle F, F^{i}\right\rangle=\left\langle F, F^{i}\right\rangle . \tag{5.4}
\end{equation*}
$$

In addition, at the point $p_{0}, \nabla_{i} \bar{W}\left(p_{0}\right)=0$ yields

$$
K^{\alpha} \nabla_{i} b^{11}=-b^{11} \nabla_{i} K^{\alpha}+\frac{n \alpha-1}{2 n \alpha} \nabla_{i}|F|^{2}=-b^{11} h_{i i}\left\langle F, F^{i}\right\rangle+\frac{n \alpha-1}{n \alpha}\left\langle F, F_{i}\right\rangle=\left(\beta-\theta_{i}\right)\left\langle F, F_{i}\right\rangle,
$$

where $\theta_{i}=b^{11} h_{i i}\left(p_{0}\right)$ and $\beta=\frac{n \alpha-1}{n \alpha}$. We also have

$$
\begin{equation*}
\left.\left.\langle F, \nabla| F\right|^{2}\right\rangle:=\left\langle F,\left(\nabla_{i}|F|^{2}\right) F^{i}\right\rangle=\left\langle F, F^{i}\right\rangle\left(\nabla_{i}|F|^{2}\right)=2\left\langle F, F_{i}\right\rangle\left\langle F, F^{i}\right\rangle . \tag{5.5}
\end{equation*}
$$

Hence, we can rewrite (5.3) as

$$
\begin{equation*}
0 \geqslant \sum_{i \neq 1}\left\langle F, F_{i}\right\rangle^{2} J_{i}+\left\langle F, F_{1}\right\rangle^{2} I_{1}+(n \alpha-1) K^{\alpha}\left(b^{11}-\frac{1}{n} \sum_{i=1}^{n} b^{i i}\right), \tag{5.6}
\end{equation*}
$$

where

$$
I_{1}=\frac{n \alpha-1}{n \alpha}+1-\alpha, \quad J_{i}=2 \alpha\left(\beta-\theta_{i}\right)+\alpha\left(\theta_{i}^{-2}+2 \theta_{i}^{-1}\right)\left(\beta-\theta_{i}\right)^{2}+\beta
$$

We observe that $I_{1}>0$ holds, and also $J_{i}$ satisfies

$$
\begin{aligned}
J_{i} & =2 \alpha \beta-2 \alpha \theta_{i}+\alpha \beta^{2} \theta_{i}^{-2}+2 \alpha \beta^{2} \theta_{i}^{-1}-2 \alpha \beta \theta_{i}^{-1}-4 \alpha \beta+\alpha+2 \alpha \theta_{i}+\beta \\
& =\alpha(1-\beta)+\beta(1-\alpha)+2 \alpha \beta(\beta-1) \theta_{i}^{-1}+\alpha \beta^{2} \theta_{i}^{-2}=\frac{1}{n}+\beta(1-\alpha)-\frac{2 \beta}{n} \theta_{i}^{-1}+\alpha \beta^{2} \theta_{i}^{-2} \\
& =\beta(1-\alpha)+\frac{1}{n}+\alpha\left(\beta \theta_{i}^{-1}-\frac{1}{n \alpha}\right)^{2}-\frac{1}{n^{2} \alpha} \geqslant \beta(1-\alpha)+\frac{1}{n}\left(\frac{n \alpha-1}{n \alpha}\right)=\beta\left(1-\alpha+\frac{1}{n}\right)>0 .
\end{aligned}
$$

Since we have $b^{11}\left(p_{0}\right)=\lambda_{\min }^{-1}\left(p_{0}\right) \geqslant \lambda_{i}^{-1}\left(p_{0}\right) \geqslant b^{i i}\left(p_{0}\right)$ and $\left\langle F, F_{i}\right\rangle^{2}\left(p_{0}\right) \geqslant 0$ for all $i \in$ $\{1, \cdots, n\}$, the inequality (5.6) and $I_{1}, J_{i}>0$ give the desired result.

## 2 Strong maximum principle

In this section, we will show how Theorem 5.5 can be modified to give us the proof of our main result, Theorem 5.1. To this end, we will introduce the new geometric, chart-independent quantity

$$
Z(p)=\left(K^{\alpha} b^{i j} g_{i j}-\frac{n \alpha-1}{2 \alpha}|F|^{2}\right)(p)
$$

and apply the strong maximum principle. If we use $W(p), h_{i j}$ can be diagonalized only at one given point. However, if we employ $Z(p)$, we can diagonalize $h_{i j}$ at each point. We begin with the following observation which simply follows from Theorem 5.5.

Proposition 5.6. Let $\Sigma$ be a strictly convex smooth closed solution of $\left({\omega *^{\alpha}}^{\alpha}\right)$ for an exponent $\alpha \in\left(\frac{1}{n}, 1+\frac{1}{n}\right)$. Assume that $F: M^{n} \rightarrow \mathbb{R}^{n+1}$ is a smooth immersion such that $F\left(M^{n}\right)=\Sigma$, and the continuous function $Z(p)$ attains its maximum at a point $p_{0}$. Then, $F\left(p_{0}\right)$ is an umbilical point and $\nabla|F|^{2}\left(p_{0}\right)=0$ holds.

Proof. We observe $b^{i j} g_{i j}(p)=\sum_{i=1}^{n} \lambda_{i}^{-1}(p)$, where $\lambda_{1}(p), \cdots, \lambda_{n}(p)$ are the principal curvatures of $\Sigma$ at $F(p)$. Therefore, we have $Z(p) \leqslant n W(p)$. However, if $W\left(p_{0}\right)=\max _{p \in M^{n}} W(p)$, then $Z\left(p_{0}\right)=n w\left(p_{0}\right)$ holds, because $F\left(p_{0}\right)$ is an umbilical point by Theorem5.5. Hence, we have

$$
Z(p) \leqslant n W(p) \leqslant \max _{p \in M^{n}} n W(p)=\max _{p \in M^{n}} Z(p)
$$

Thus, if $Z$ attains its maximum at a point $p_{0}$, then $W$ also attains its maximum at $p_{0}$, and thus we can obtain the desired result by Theorem 5.5 .

We will now employ the strong maximum principle to prove Theorem 5.1 .
Proof of Theorem 5.1. We define a set $M_{Z} \subset M^{n}$ by

$$
M_{Z}=\left\{p \in M^{n}: Z(p)=\max _{M^{n}} Z\right\}
$$

Since $Z(p)$ is a continuous function defined on a closed manifold $M^{n}, Z$ attains its maximum, and thus $M_{Z}$ is not an empty set. We now define the continuous function $\Lambda: M^{n} \rightarrow \mathbb{R}$ and the open set $V \subset M^{n}$ by

$$
\Lambda(p)=\sum_{i, j}\left(\frac{\lambda_{i}}{\lambda_{j}}-\frac{\lambda_{j}}{\lambda_{i}}\right)^{2}(p), \quad V=\left\{p \in M^{n}: \Lambda(p)<\left(\frac{10}{9}-\frac{9}{10}\right)^{2}\right\}
$$

We now begin by combining (2.17) and (2.18) to obtain

$$
\begin{aligned}
\mathcal{L}\left(K^{\alpha} b^{p q}\right)= & 2\left\langle\nabla K^{\alpha}, \nabla b^{p q}\right\rangle_{\mathcal{L}}+b^{p r} b^{q s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}+\alpha K^{2 \alpha} b^{p r} b^{q s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m} \\
& +\left\langle F, \nabla\left(K^{\alpha} b^{p q}\right)\right\rangle+(n \alpha-1) K^{\alpha}\left(b^{p q}-g^{p q} K^{\alpha}\right) .
\end{aligned}
$$

Therefore, we can derive the following from (2.15) and $\nabla g_{p q}=0$

$$
\mathcal{L} Z=2 g_{p q}\left\langle\nabla K^{\alpha}, \nabla b^{p q}\right\rangle_{\mathcal{L}}+b^{p r} b_{p}^{s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}+\alpha K^{2 \alpha} b^{p r} b_{p}^{s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m}+\left\langle F, \nabla\left(K^{\alpha} b^{p q} g_{p q}\right)\right\rangle .
$$

By using (5.5), we can obtain

$$
\left.\left\langle F, \nabla\left(K^{\alpha} b^{p q} g_{p q}\right)\right\rangle=\langle F, \nabla Z\rangle+\left.\frac{n \alpha-1}{2 \alpha}\langle F, \nabla| F\right|^{2}\right\rangle=\langle F, \nabla Z\rangle+\left(n-\alpha^{-1}\right)\left\langle F, F_{i}\right\rangle\left\langle F, F^{i}\right\rangle .
$$

Hence, we have

$$
\begin{align*}
\mathcal{L} Z-\langle F, \nabla Z\rangle= & 2 \alpha\left(b^{i j} \nabla_{i} K^{\alpha}\right)\left(K^{\alpha} g_{p q} \nabla_{j} b^{p q}\right)+b^{p r} b_{p}^{s} \nabla_{r} K^{\alpha} \nabla_{s} K^{\alpha}  \tag{5.7}\\
& +\alpha K^{2 \alpha} b^{p r} b_{p}^{s} b^{i j} b^{k m} \nabla_{r} h_{i k} \nabla_{s} h_{j m}+\left(n-\alpha^{-1}\right)\left\langle F, F_{i}\right\rangle\left\langle F, F^{i}\right\rangle .
\end{align*}
$$

Given a fixed point $p_{0} \in V$, we choose an orthonormal frame at $F\left(p_{0}\right)$ satisfying

$$
g_{i j}\left(p_{0}\right)=\delta_{i j}, \quad h_{i j}\left(p_{0}\right)=\lambda_{i}\left(p_{0}\right) \delta_{i j} .
$$

Then, at the point $p_{0}$, we can rewrite (5.7) as

$$
\begin{align*}
\mathcal{L} Z-\langle F, \nabla Z\rangle= & 2 \alpha \sum_{i, j}\left(b^{i i} \nabla_{i} K^{\alpha}\right)\left(K^{\alpha} \nabla_{i} b^{j j}\right)+\sum_{i}\left|b^{i i} \nabla_{i} K^{\alpha}\right|^{2}  \tag{5.8}\\
& +\alpha K^{2 \alpha} \sum_{i, j, k}\left(b^{i i}\right)^{2} b^{j j} b^{k k}\left|\nabla_{i} h_{j k}\right|^{2}+\left(n-\alpha^{-1}\right) \sum_{i}\left\langle F, F_{i}\right\rangle^{2} .
\end{align*}
$$

Since $p_{0} \in V$ and the definition of $V$ guarantees that $b^{i i} h_{j j}\left(p_{0}\right) \geqslant \frac{9}{10}$, by using (2.14) we can derive

$$
\begin{aligned}
\alpha K^{2 \alpha} \sum_{i, j, k}\left(b^{i i}\right)^{2} b^{j j} b^{k k}\left|\nabla_{i} h_{j k}\right|^{2} & \geqslant \alpha \sum_{i}\left|K^{\alpha} \nabla_{i} b^{i i}\right|^{2}+2 \alpha \sum_{i \neq j} b^{j j} h_{i i}\left|K^{\alpha} \nabla_{j} b^{i i}\right|^{2}+\alpha \sum_{i \neq j}\left(b^{i i} h_{j j}\right)^{2}\left|K^{\alpha} \nabla_{i} b^{j j}\right|^{2} \\
& \geqslant \alpha \sum_{i}\left|K^{\alpha} \nabla_{i} b^{i i}\right|^{2}+\frac{5}{2} \alpha \sum_{i \neq j}\left|K^{\alpha} \nabla_{i} b^{j j}\right|^{2}
\end{aligned}
$$

We also have

$$
\alpha \sum_{i}\left|K^{\alpha} \nabla_{i} b^{i i}\right|^{2}+2 \alpha \sum_{i}\left(b^{i i} \nabla_{i} K^{\alpha}\right)\left(K^{\alpha} \nabla_{i} b^{i i}\right) \geqslant-\alpha \sum_{i}\left|b^{i i} \nabla_{i} K^{\alpha}\right|^{2}
$$

and
$\frac{5}{2} \alpha \sum_{i \neq j}\left|K^{\alpha} \nabla_{i} b^{j j}\right|^{2}+2 \alpha \sum_{i \neq j}\left(b^{i i} \nabla_{i} K^{\alpha}\right)\left(K^{\alpha} \nabla_{i} b^{j j}\right) \geqslant-\frac{2}{5} \alpha \sum_{i \neq j}\left|b^{i i} \nabla_{i} K^{\alpha}\right|^{2}=-\frac{2}{5} \alpha(n-1) \sum_{i}\left|b^{i i} \nabla_{i} K^{\alpha}\right|^{2}$.
Applying the inequalities above and (5.4) to (5.8) yields

$$
\begin{aligned}
\mathcal{L} Z-\langle F, \nabla Z\rangle & \geqslant\left((1-\alpha)-\frac{2}{5} \alpha(n-1)+\left(n-\alpha^{-1}\right)\right) \sum_{i}\left\langle F, F_{i}\right\rangle^{2} \\
& =\frac{1}{5 \alpha}\left(-(2 n+3) \alpha^{2}+5(n+1) \alpha-5\right) \sum_{i}\left\langle F, F_{i}\right\rangle^{2} .
\end{aligned}
$$

Let us consider the function $y(\alpha)=-(2 n+3) \alpha^{2}+5(n+1) \alpha-5$. Then, we have

$$
y(1+1 / n)=3 n-2-(3 / n)-\left(3 / n^{2}\right) \geqslant 0, \quad y(1 / n)=(3 / n)-\left(3 / n^{2}\right) \geqslant 0,
$$

which implies $y(\alpha) \geqslant 0$ for $\alpha \in\left[\frac{1}{n}, 1+\frac{1}{n}\right]$. Therefore, on $V$ the following holds

$$
\mathcal{L} Z-\langle F, \nabla Z\rangle \geqslant 0 .
$$

Notice that $\mathcal{L} Z-\langle F, \nabla Z\rangle$ is a chart-independent function. Hence, the Hopf maximum principle and $M_{Z} \subset V$ show that $M_{Z}=V$. However, $M_{Z}$ is a closed set and $V$ is an open set by the continuity of $Z$ and $\Lambda$, respectively. So, we conclude that $M_{Z}=M^{n}$, and thus Proposition 5.6 gives the desired result.

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